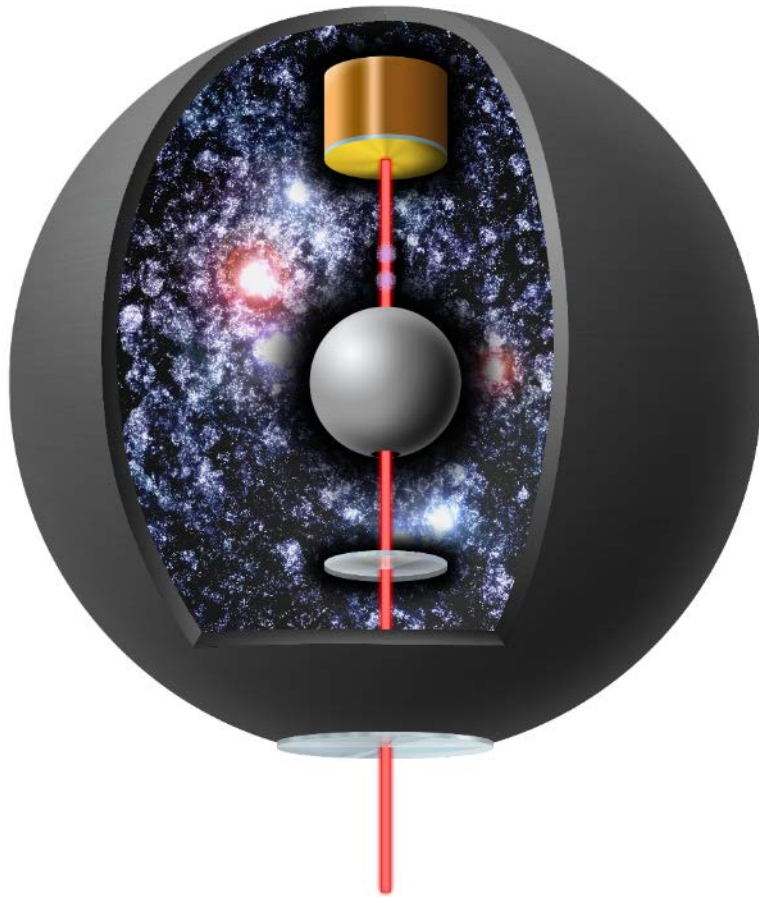




UCLA



Fundamental physics with matter wave interferometry*

Paul Hamilton

University of California at
Los Angeles

W. M. KECK FOUNDATION



UNIVERSITY OF CALIFORNIA Office of the President



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Outline

1. Introduction to matter wave interferometry
2. Berkeley dark energy search
3. UCLA Bloch oscillation accelerometer
4. HUNTER – sterile neutrino search

de Broglie waves

It all starts with a Ph.D. thesis!
(de Broglie, 1924)

$$E = mc^2 = \hbar\omega_c$$

Lorentz invariance requires

$$\Psi \sim \exp(-i\omega_c\tau)$$

- Holds in all of QM and QFT

$$\lambda_{dB} = \frac{h}{mv}$$

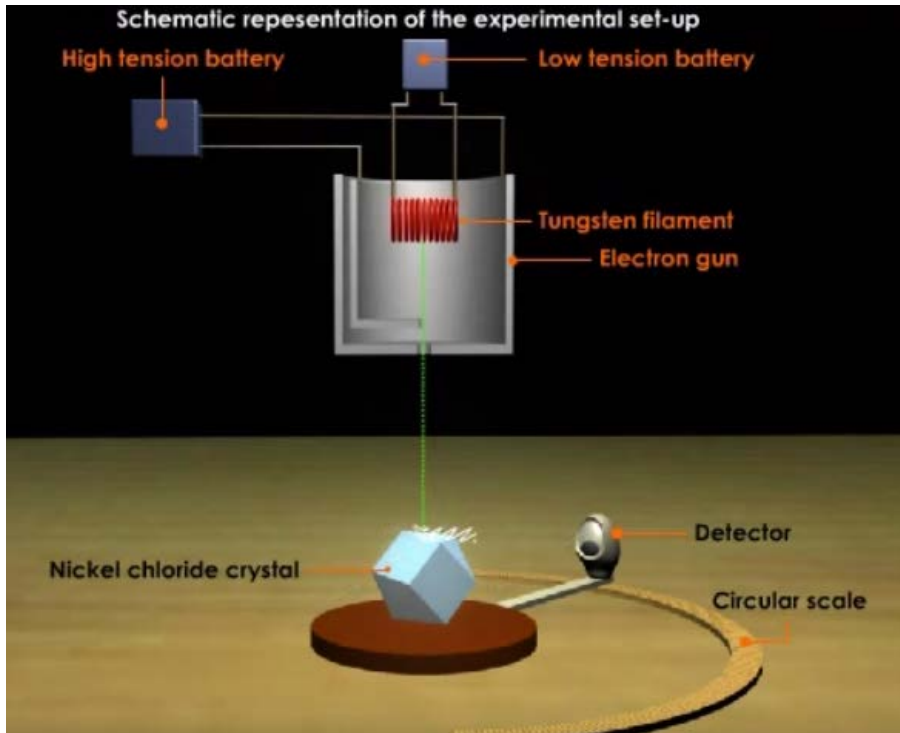




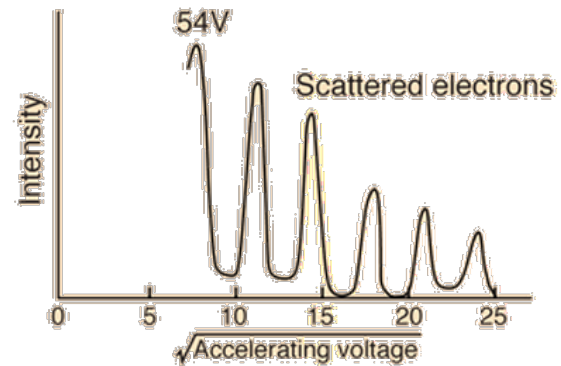
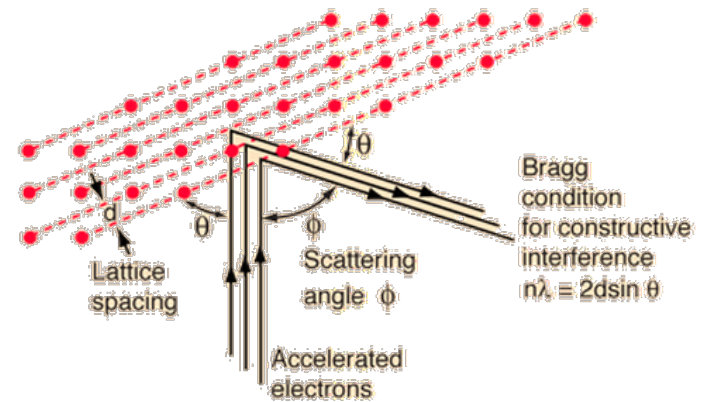
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From electrons...

1927 Davisson Germer experiment

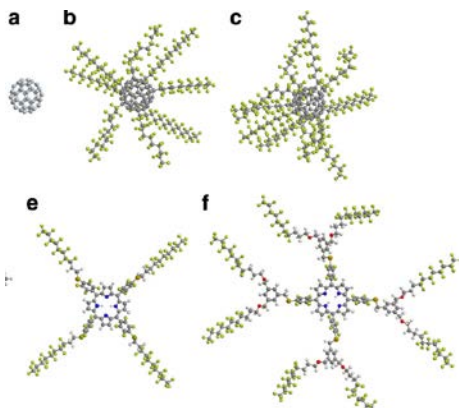
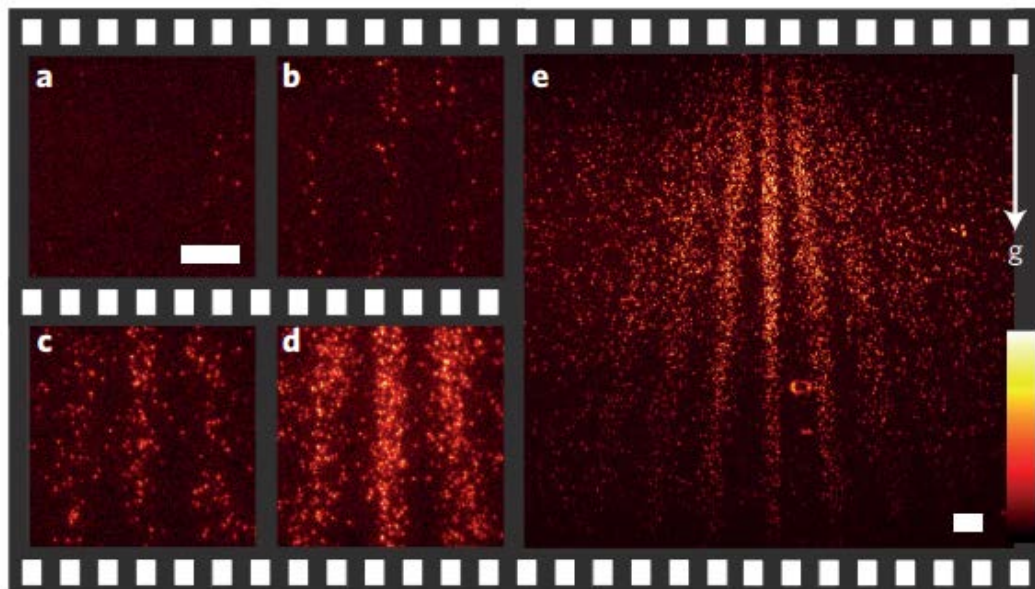
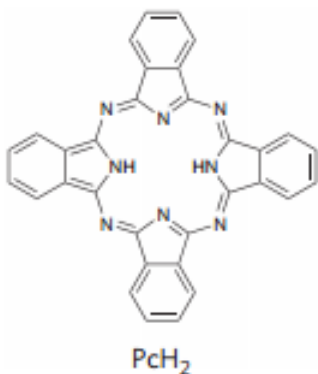


Credit: Tutorvista.com



Davisson, C. J., "Are Electrons Waves?," Franklin Institute Journal 205, 597 (1928)

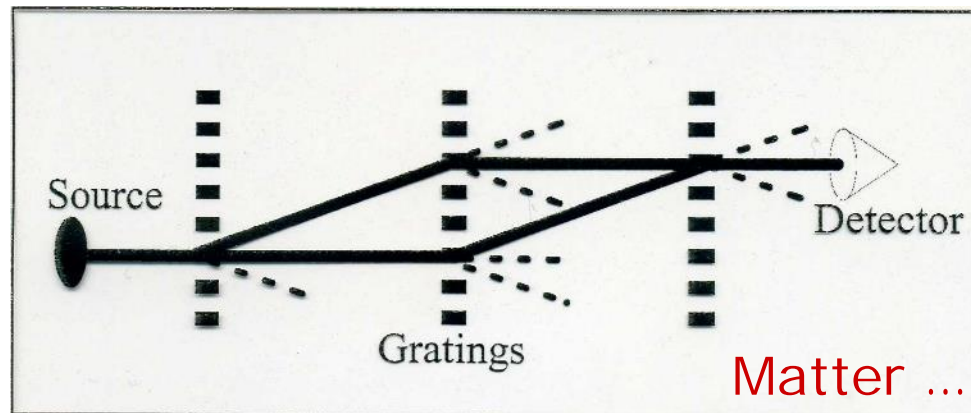
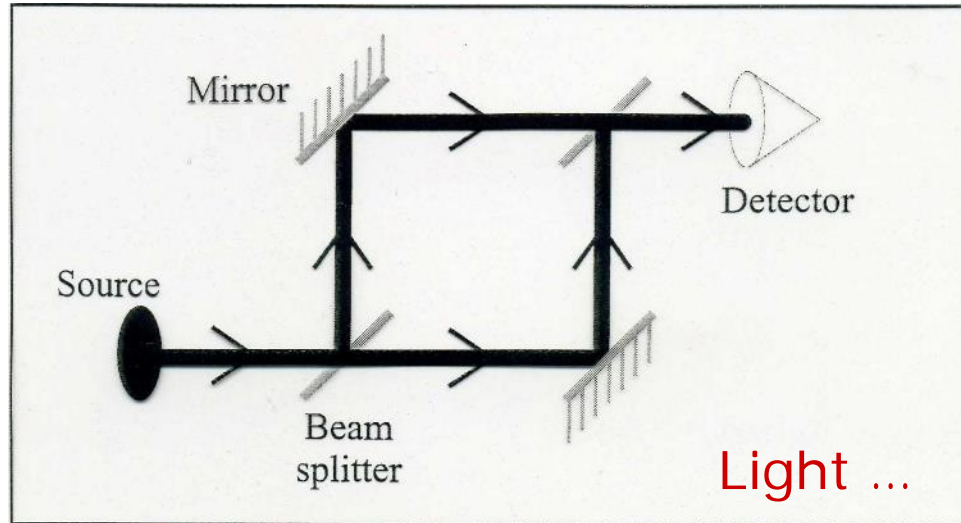
...to large molecules



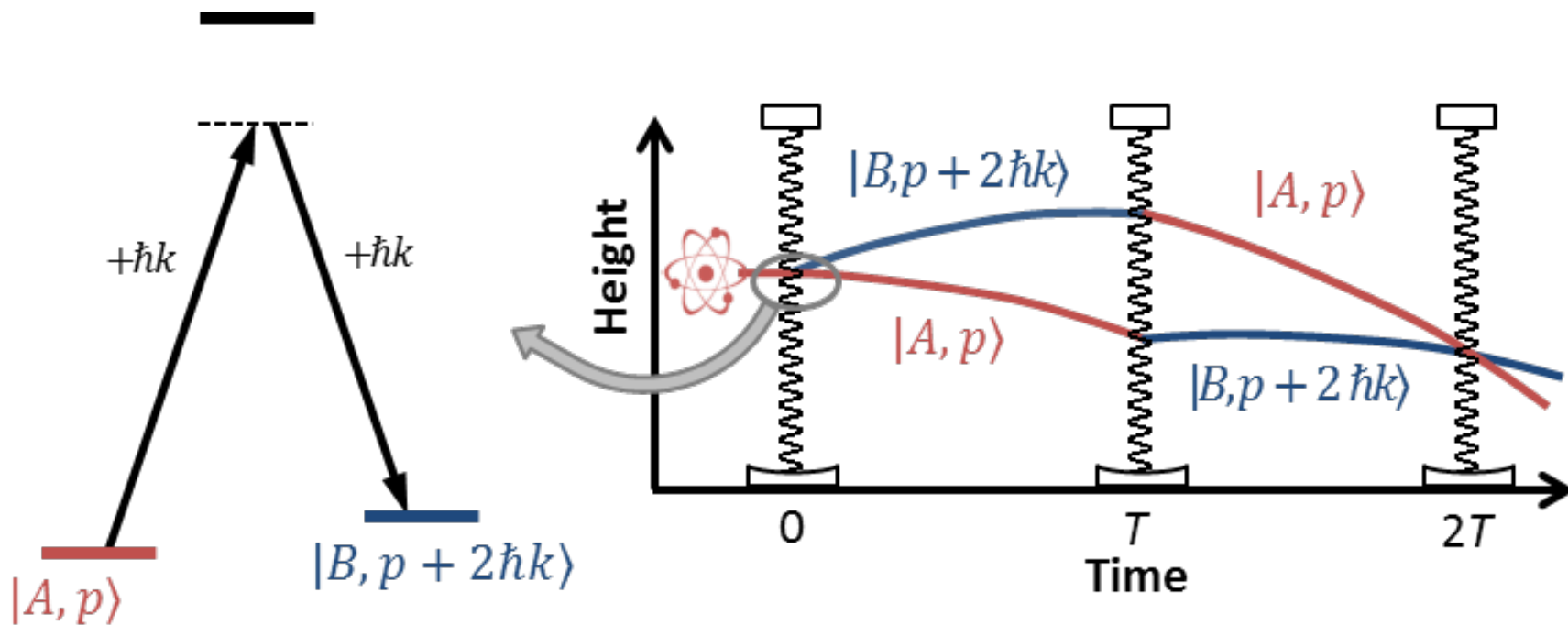
Buckyballs and molecules with thousands of atoms interfered by Zeilinger & Arndt groups in Vienna

Working towards viruses and bacteria

Modern matter interferometry



Light pulse atom interferometer



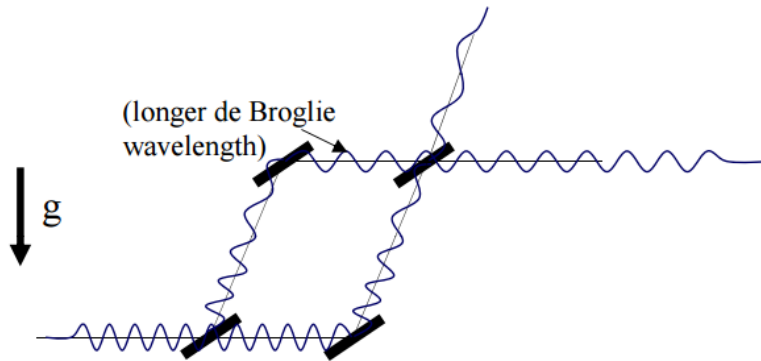
$$\Delta\varphi = -\frac{1}{\hbar} \oint L d\tau + \Delta\varphi_{\text{laser}}$$

$$= \vec{k} \cdot \vec{g} T^2$$

Measuring all accelerations

Gravity/Accelerations

As atom climbs gravitational potential, velocity decreases and wavelength increases



Demonstrated sensitivity of atom interferometers

- Accelerations: $< \text{ppb of } g!$
- Rotations: $< \text{nrad/s !}$

Rotations

Sagnac effect for de Broglie waves

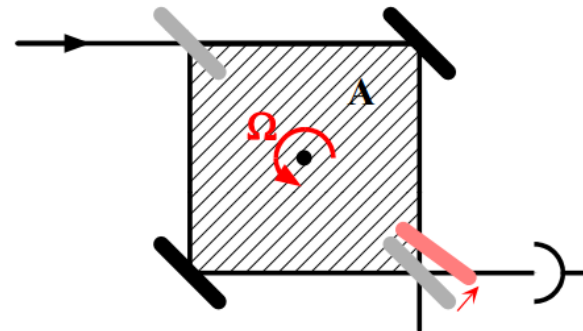


Figure credit: Kasevich

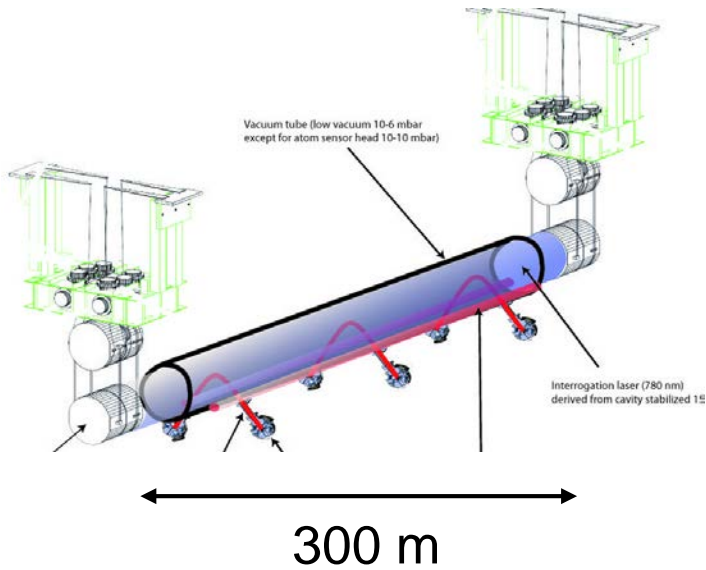
Field uses:

- Inertial navigation
- Mineral and oil searches
- Hydrology
- Proposed for geodesy



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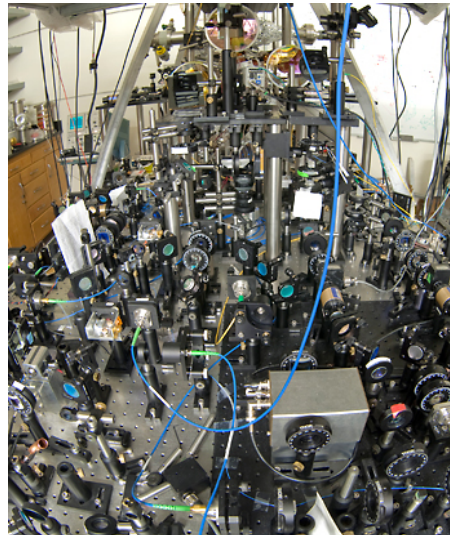
Applications



Gravity wave detection
MIGA collaboration



Tests of general relativity
Stanford 10 meter atomic fountain



Fine structure constant
Measurement
3m atomic fountain
Berkeley



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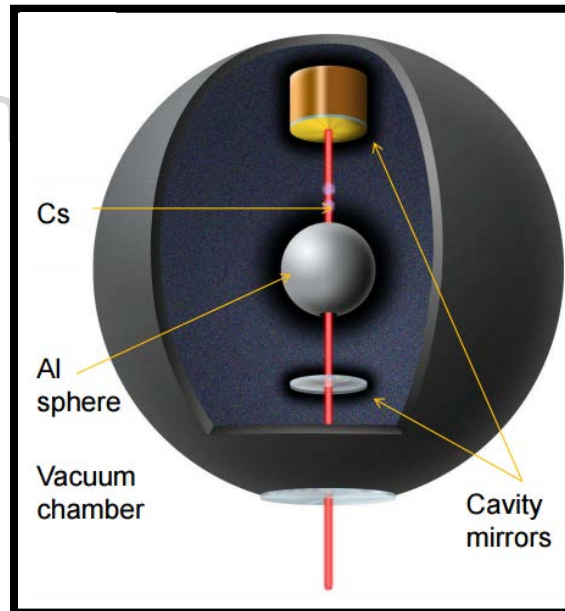
Outline

1. Intro to matter wave interferometry

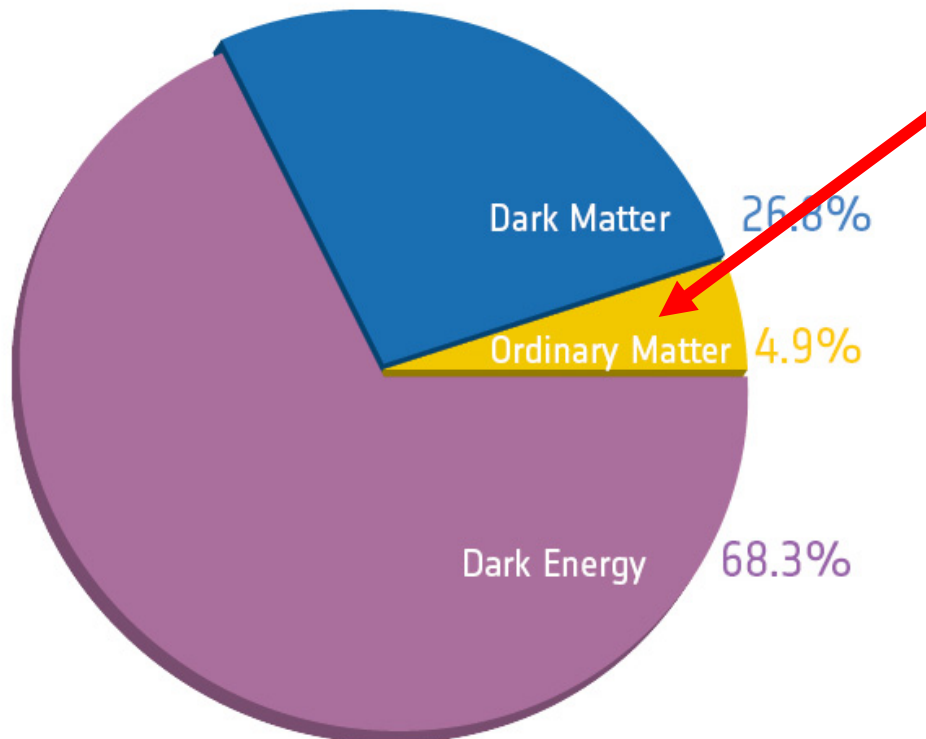
2. Berkeley dark energy search

3. UCLA Bloch interferometer

4. HUNTER – dark energy search



Contents of the Universe



ESA/Planck

We know this part very well!!

Decades of testing and no definitive violations of SM.

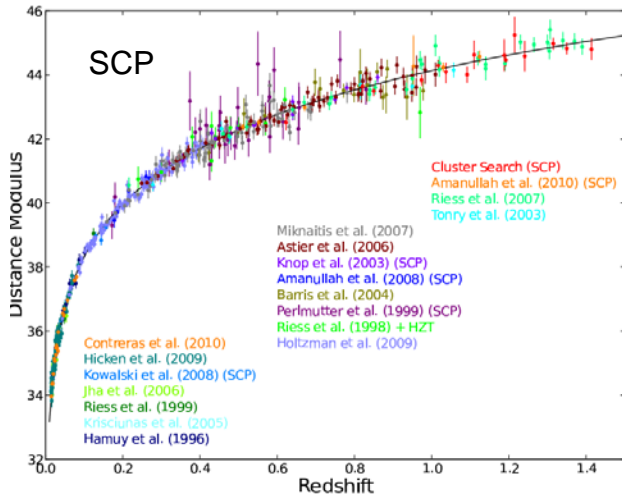
But what is the rest?

What is dark energy?

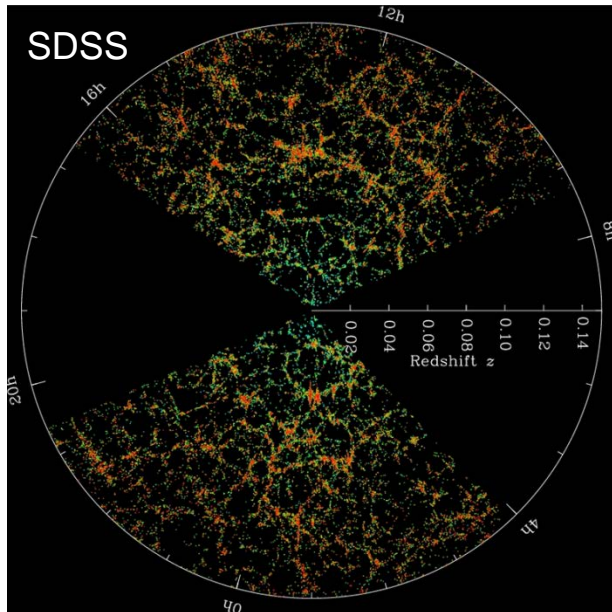
What is dark matter?

Why matter and not antimatter?

What can astronomers learn?



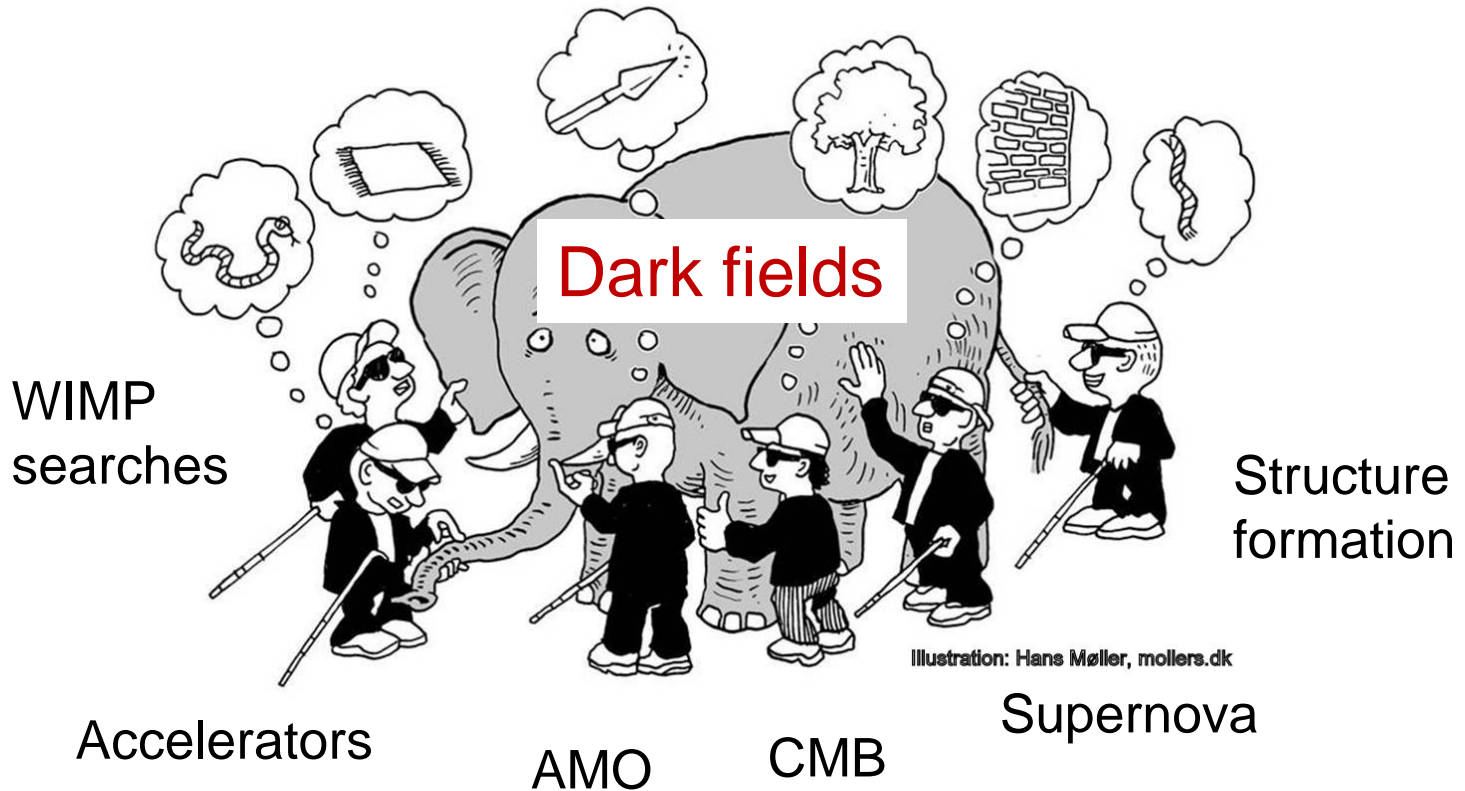
- Dark energy density, Λ hydrogen / m^3 or an energy scale of 2.4 meV
- Equation of state for dark energy
- Time evolution of dark energy?



Very large-scale experiments for measuring single properties of dark energy.

In the dark...

AMO observations can help
test complementary
properties



How can we detect scalar fields in the laboratory?

Measuring local expansion of the universe is beyond reach.

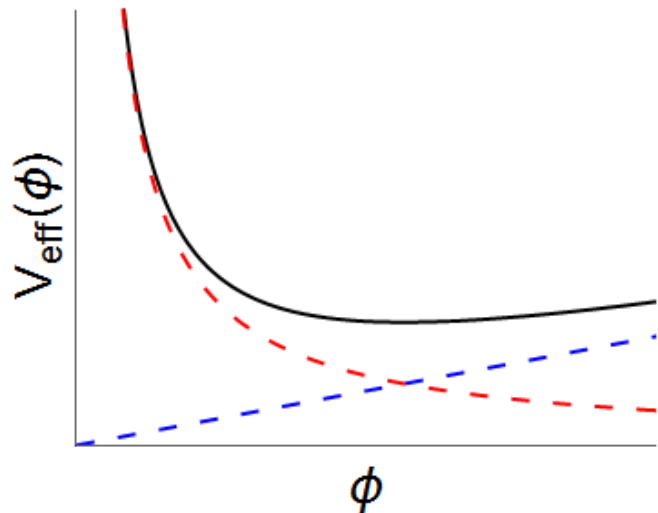
...but **new fields** can lead to **new forces** between objects.

Quintessence + coupling

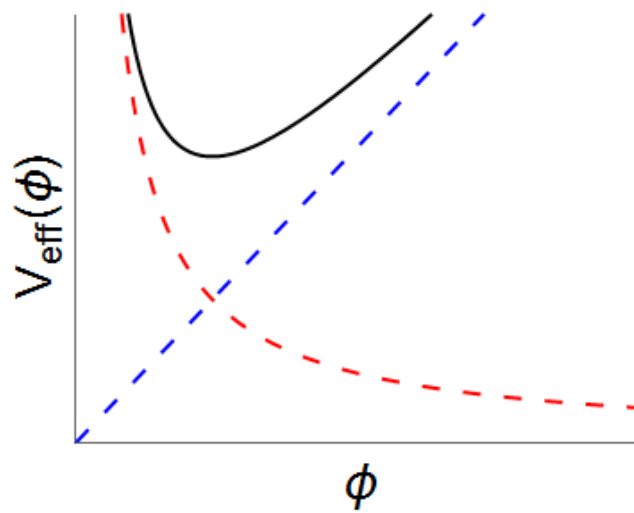
Quintessence: dark energy = scalar field (Ratra, Peebles 1988!)

$$V_{\text{eff}} = \Lambda^4 + \frac{\Lambda^{4+n}}{\phi^n} + \frac{\phi}{M} \rho$$

“Textbook”
quintessence
Coupling to
local density



Low density ρ
vacuum



High density ρ
normal matter

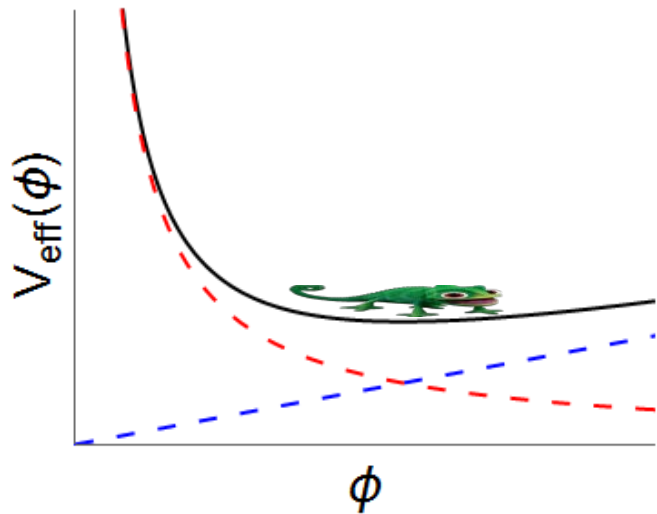
2 parameters:
 Λ and M

Chameleon mechanism

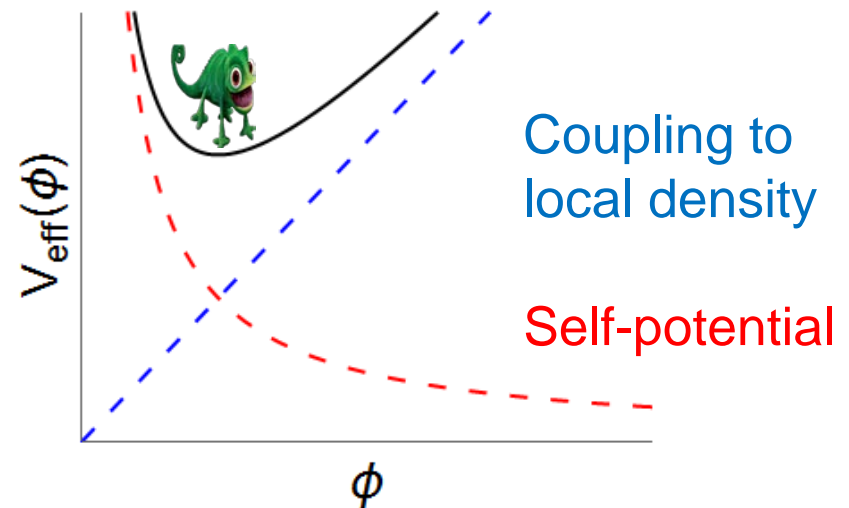


Khoury, Weltman
Phys. Rev. D **69**, 044026

Curvature \Rightarrow Mass \Rightarrow $1 / \text{range}$



In vacuum \Rightarrow long range



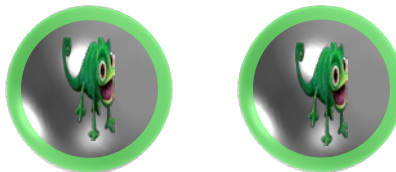
Near normal matter \Rightarrow short range

Screened force

Newtonian
force

New chameleon
force

$$F_{chameleon} = \frac{GM_A M_B}{r^2} \left[1 + 2 \lambda_A \lambda_B \left(\frac{M_{Pl}}{M} \right)^2 \right]$$



$$\lambda_i(\Lambda, M, \rho_i, r_i) = \frac{\textit{Shell mass}}{\textit{Test mass}}$$

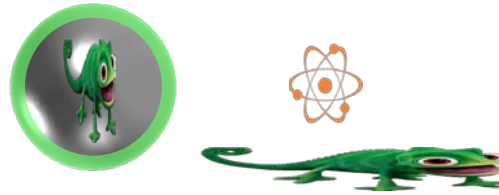
Can be extremely small ($\ll 10^{-20}$) for
macroscopic objects

$$M < M_{Pl}$$

Unscreened force can be much stronger
than gravity

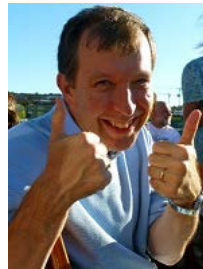
Atoms evade screening

$$F_{chameleon} = \frac{GM_A M_B}{r^2} \left[1 + 2 \lambda_A \lambda_B \left(\frac{M_{Pl}}{M} \right)^2 \right]$$



$\lambda_{atom} = 1$ For most of parameter space

Burrage, Copeland, Hinds JCAP03(2015)042

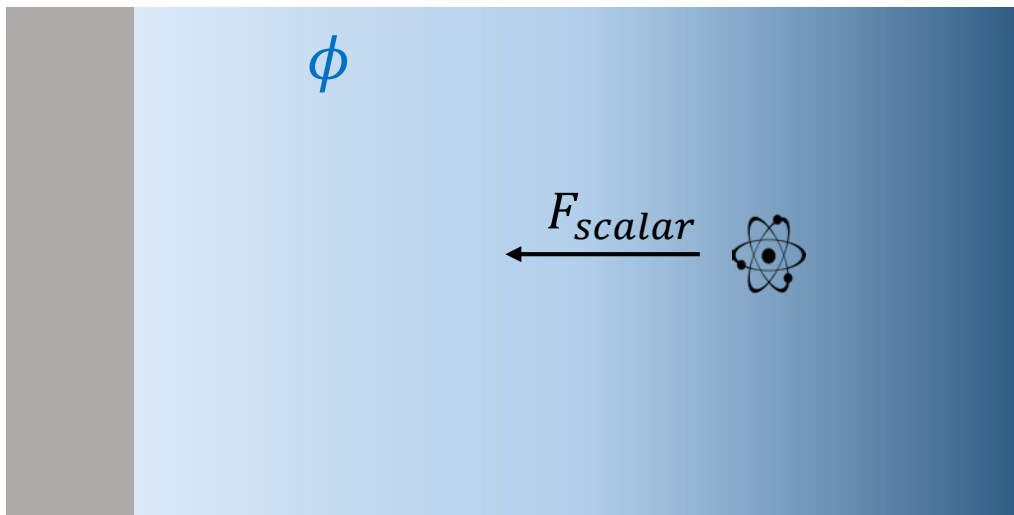


Generic predictions of forces

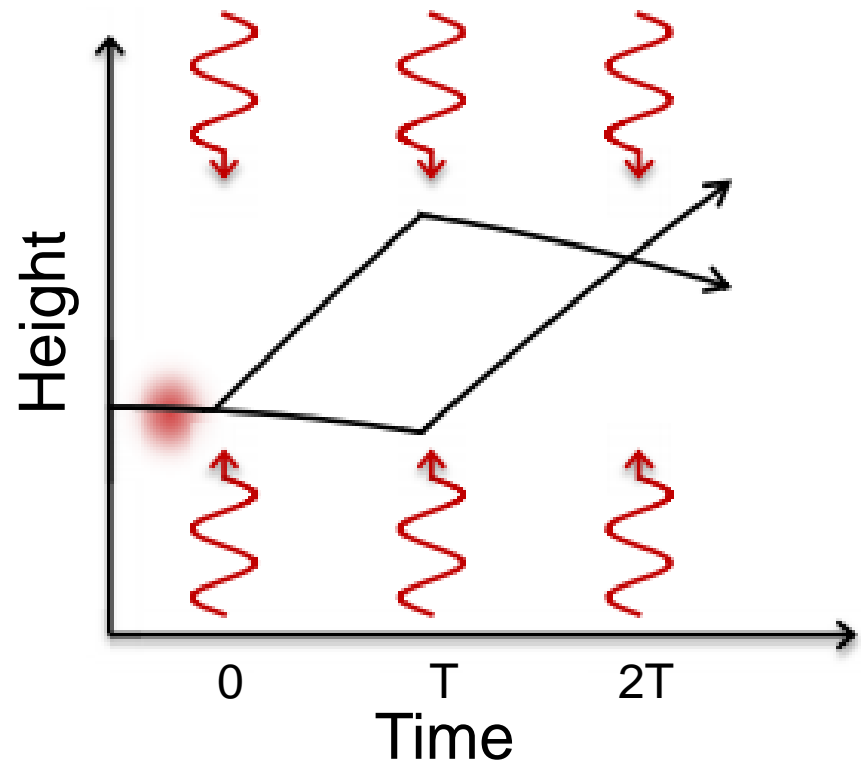
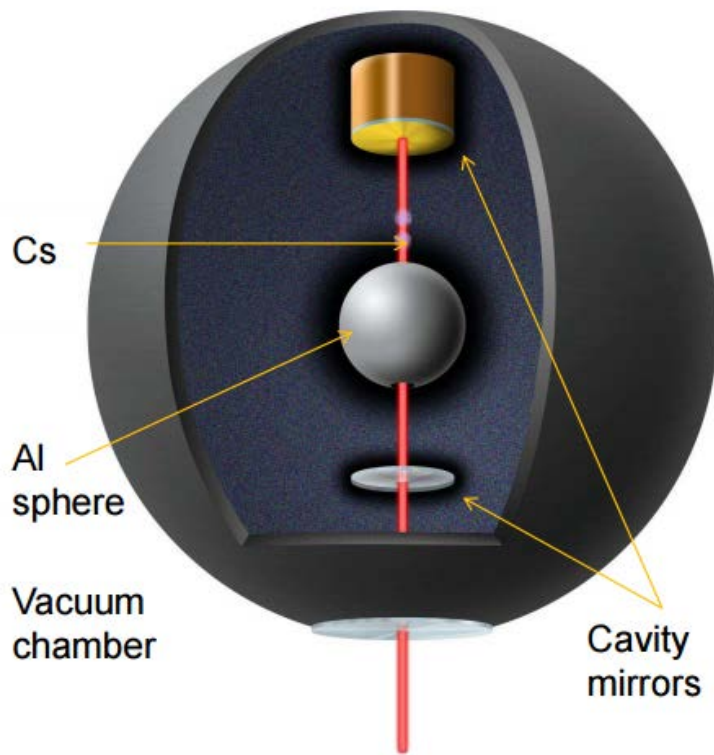
Chameleons are one example, but in general scalar fields with couplings to matter can create screened forces:

$$a_{scalar} = \frac{GM}{r^2} - \frac{\beta(\phi)}{M_{pl}} \lambda_a \nabla \phi$$

Coupling \downarrow
 Scalar gradient \swarrow
 Screening \nwarrow



Berkeley dark energy search

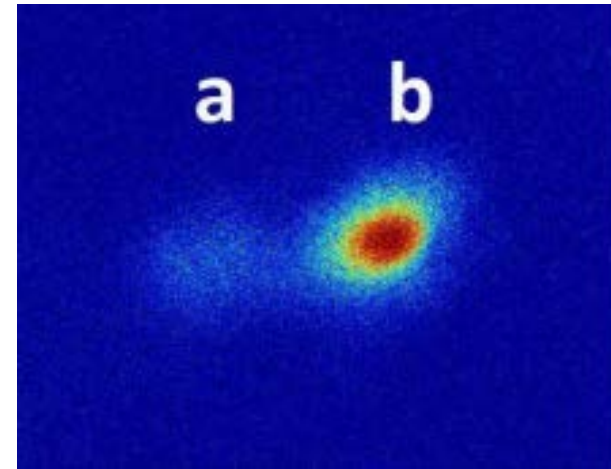
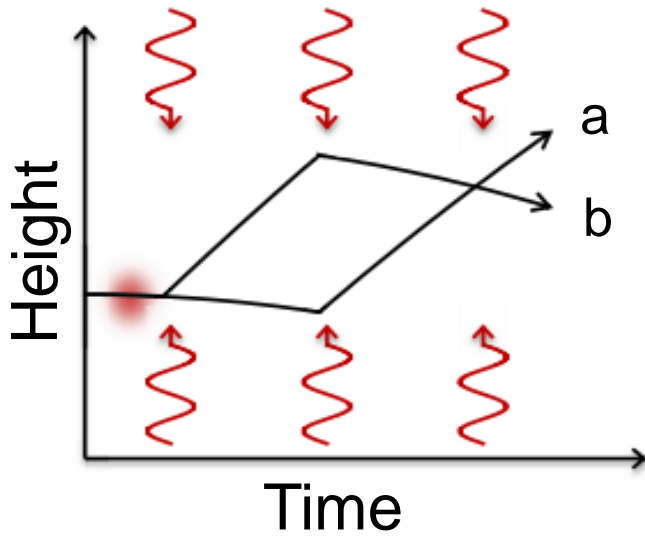


- Metal sphere creates gradient in scalar field
- Atoms act as test masses for force sensing



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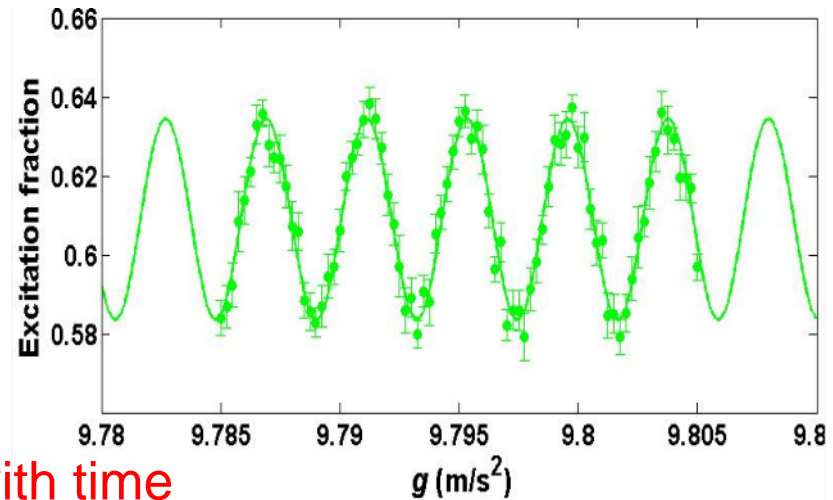
Detection

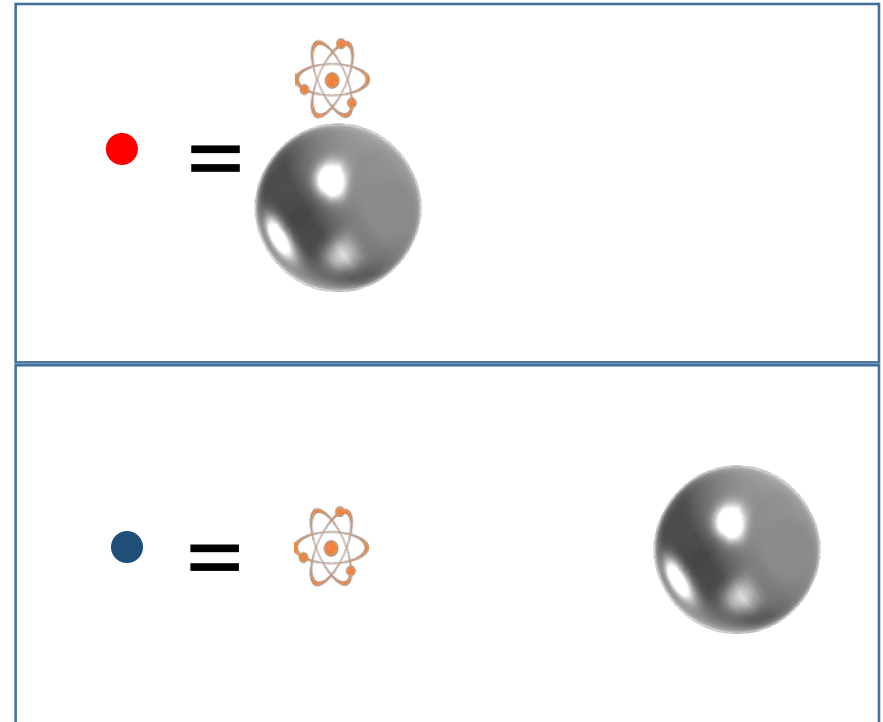
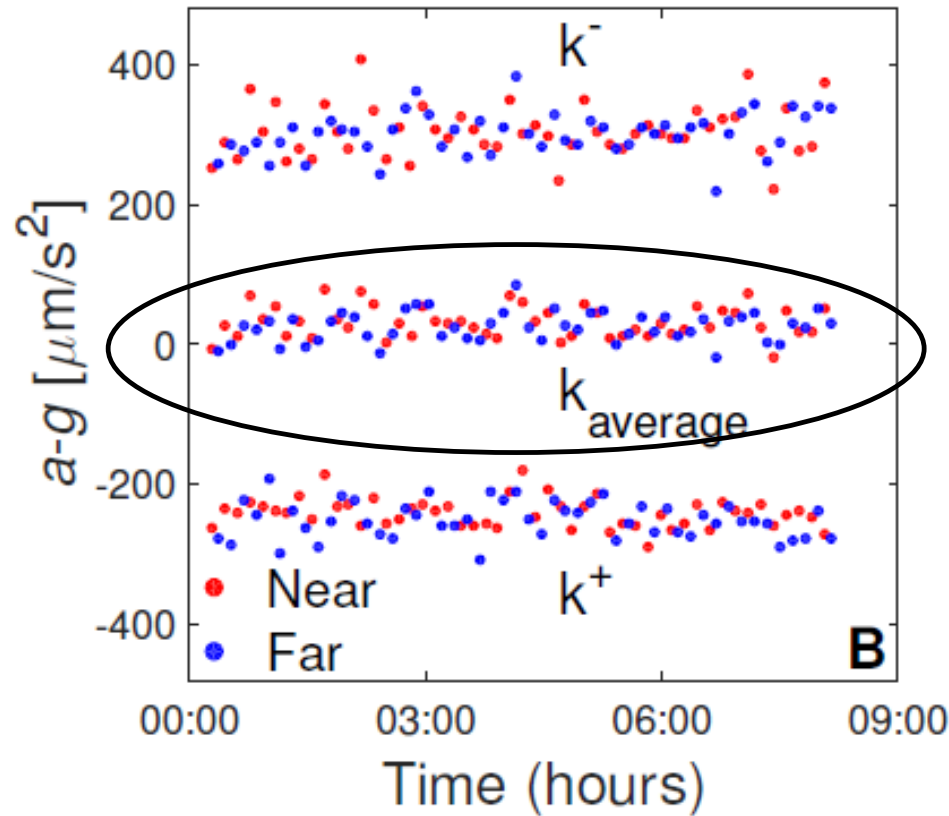


Probability
 $\propto \text{Cos}^2(\vec{k} \cdot \vec{a} T^2)$

Sensitive to
 accelerations
 → force sensor

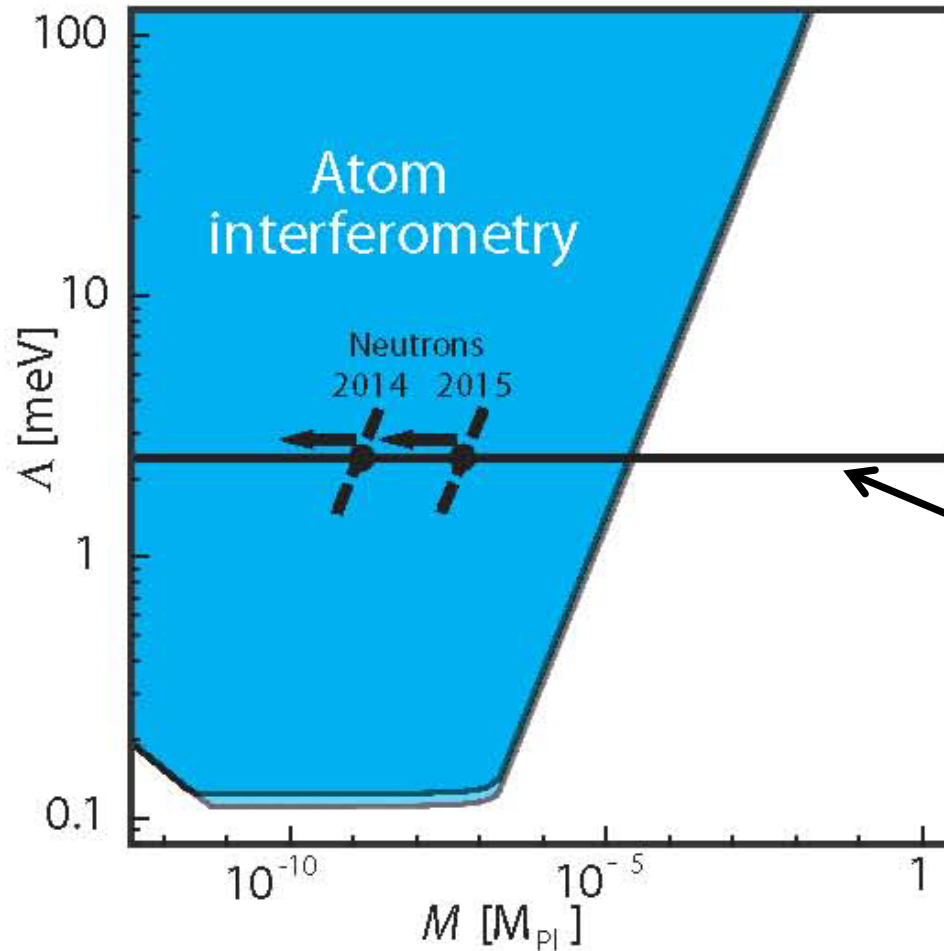
Scales with time
 near source mass





Search for an anomalous acceleration
when atoms are near the source

Dark energy limits



$$V_{\text{eff}} = \Lambda^4 + \frac{\Lambda^{4+n}}{\phi^n} + \frac{\phi}{M} \rho$$

↑ Quintessence potential
↑ Coupling to local density

Cosmological dark energy
 $\Lambda = 2.4 \text{ meV}$

$$M > 2.3 \times 10^{-5} M_{Pl}$$

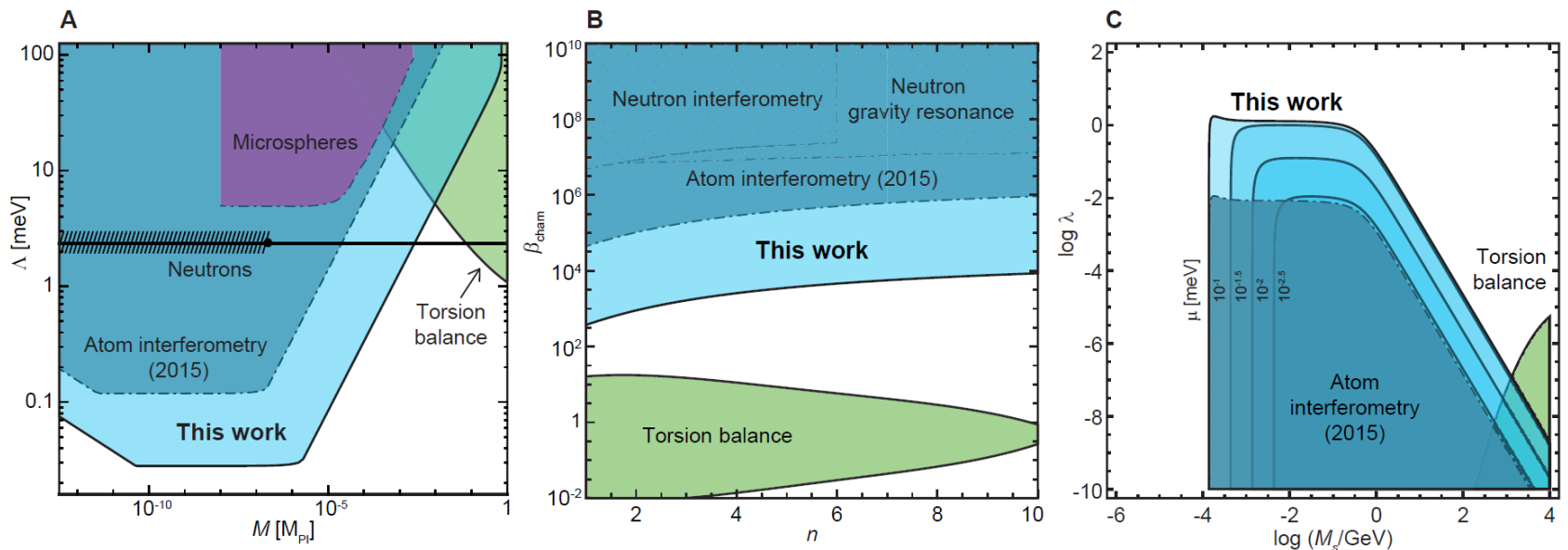
$$\Delta a = 2.3 \pm 3.3 \mu\text{m/s}^2$$

Latest results

After a year of hard work and improvements

→ $a_{\text{anomaly}} < 45 \text{ nm/s}^2$ (95% confidence)

100x improvement on chameleon and symmetron bounds

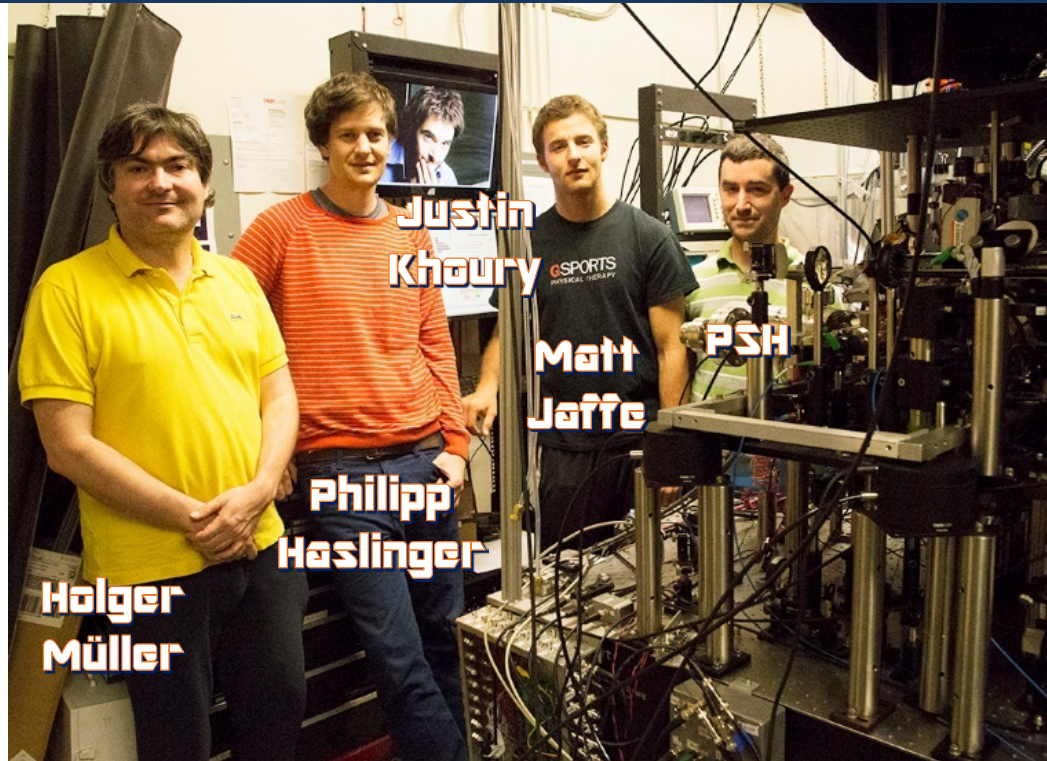


Jaffe et al., Nature Physics (2017)

Take home message: a few orders of magnitude more will either discover or rule out these theories



Acknowledgements



Dark energy search

Matt Jaffe Philipp Haslinger Victoria Xu

Benjamin Elder Justin Khoury Amol Upadhye

PI : Holger Müller

How do we go further?

The small size of the Berkeley experiment limits its absolute sensitivity:

Berkeley $\sim 10^{-6}$ g / $\sqrt{\text{Hz}}$

Gravimeters $\sim 10^{-9}$ g / $\sqrt{\text{Hz}}$

Stanford 10m fountain $\sim 10^{-12}$ g

...other experiments take advantage of free fall distances at the meter scale.

Two possibilities:

- Trapped atom interferometry
- Microgravity such as CAL / BEC CAL





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1. Intro to matter wave interferometry
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Simple CW atom interferometer



“Ideal” atom interferometer:

- Simple
- Compact
- High sensitivity
- Continuous measurement

Goal: Turn on a laser and plug the output of a detector into an oscilloscope.

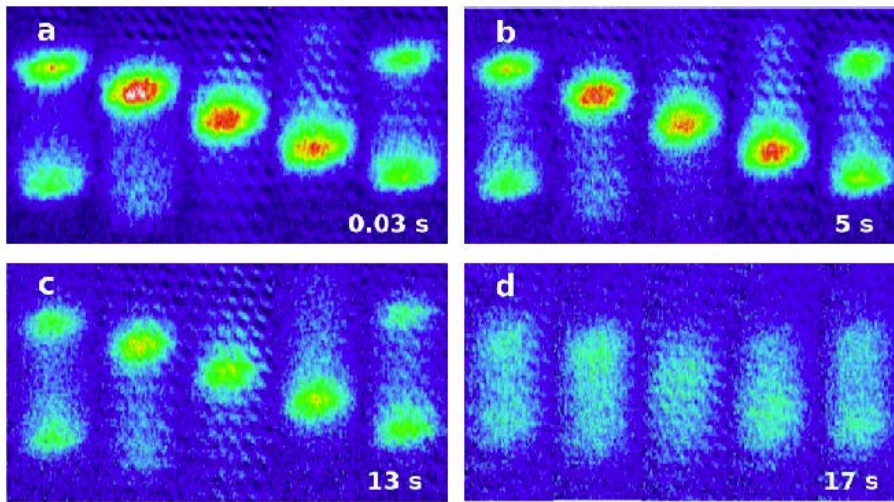
Enable measurement of AC signals

Principle: **Monitor atoms effect on a standing wave in an optical cavity**

Bloch oscillations

In quantum mechanics a force on a particle in a periodic potential leads to changes in momentum called Bloch oscillations.

$$\omega_{Bloch} = \left(F \times \frac{\lambda}{2} \right) / \hbar \quad (\sim \text{kHz scale})$$

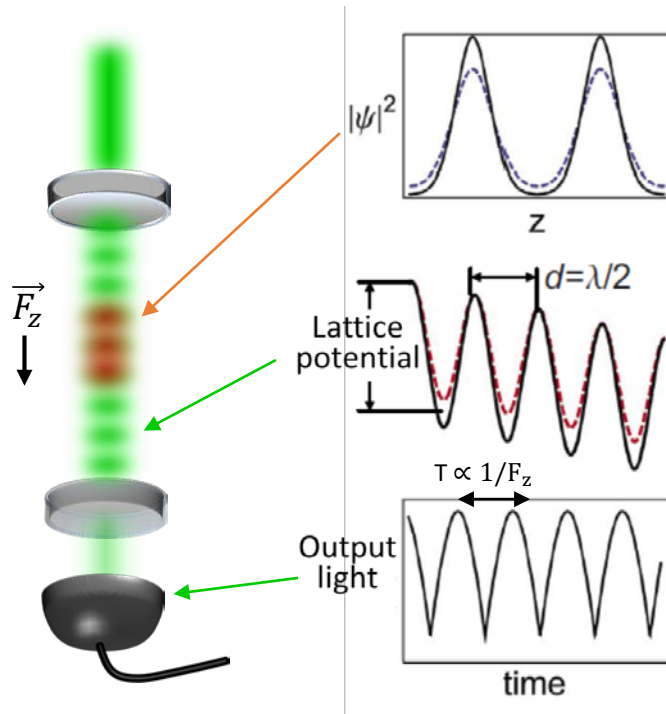


Usual method:

1. Bloch oscillations in lattice
2. Release atoms
3. Destructively image

Continuous trapped accelerometer

Collectively couple atoms to the optical cavity.



Adapted from Peden et al.

Atomic wavefunction modulates at Bloch frequency...

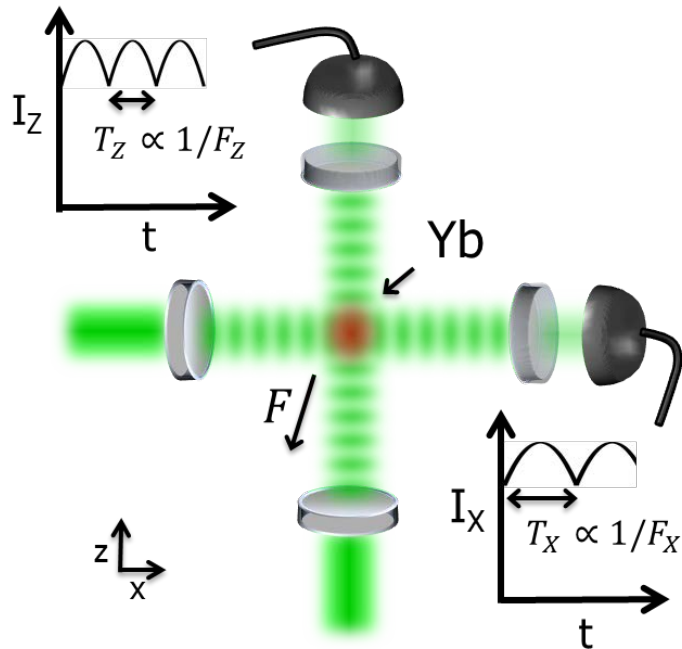
which couples to the intracavity lattice...

leading to modulation of the output light field.



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Continuous trapped accelerometer



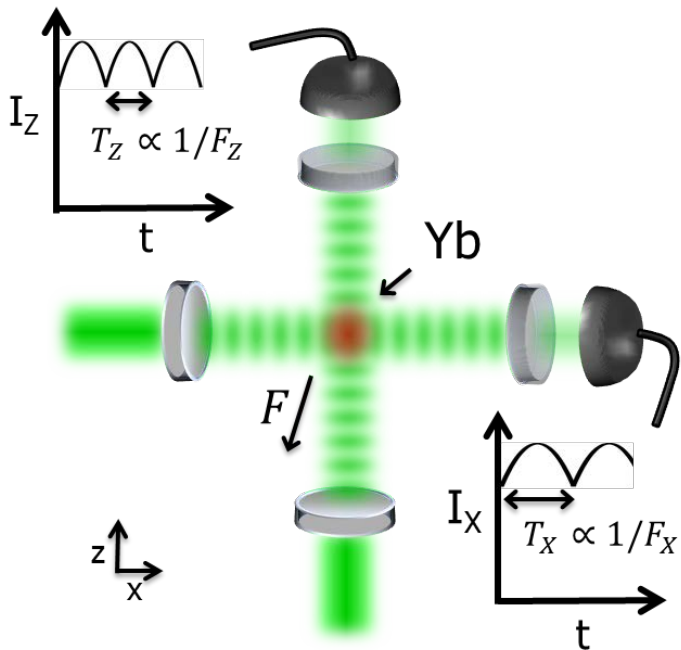
The output light of an optical cavity used to generate the optical lattice will modulate at the Bloch frequency

$$\omega_{Bloch} = \left(F \times \frac{\lambda}{2} \right) / \hbar$$

Advantages:

- Long coherence time
- Continuous readout
- Reduced vibration sensitivity
- Efficient detection

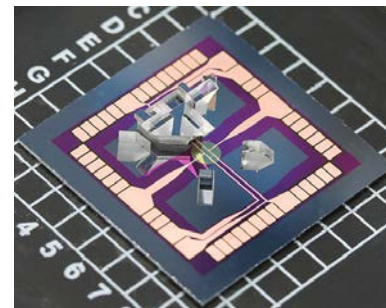
Applications



Inertial sensing:

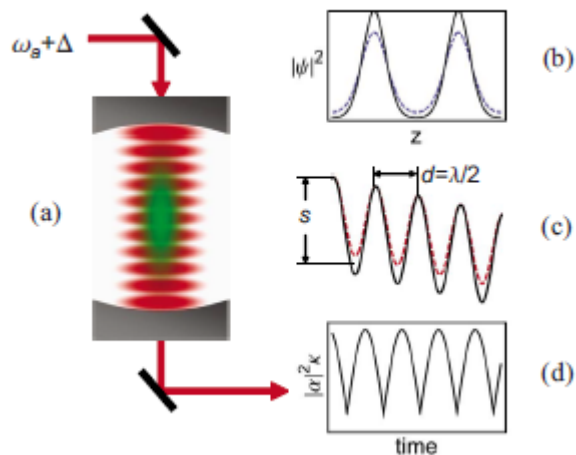
- 10^6 trapped Yb atoms
 - 5 s coherence
- $\Rightarrow \frac{10^{-8} g}{\sqrt{\text{Hz}}}$ shot-noise sensitivity

Dream sensor: integrated optical cavity/atom chip



Dana Anderson, JILA

Cavity Bloch theory



Simplified Hamiltonian

$$H = \frac{\hat{p}^2}{2M} + Mgz + \frac{\hbar g_0^2 \cos^2(k_c z) \hat{a}^\dagger \hat{a}}{\delta} - i\hbar\eta (\hat{a} - \hat{a}^\dagger) - i\hbar\kappa \hat{a}^\dagger \hat{a} .$$

Equations of motion

$$i\hbar \dot{\Psi} = \left(-\frac{\hbar^2}{2M} \partial_z^2 + \frac{\hbar g_0^2 \alpha^* \alpha}{\delta} \cos^2(k_c z) + Fz \right) \Psi$$

$$\dot{\alpha} = -i \frac{\alpha}{\delta} g^2(t) + \eta - \kappa \alpha ,$$

Peden et al. PRA 80, 043803 (2009)

α = field amplitude

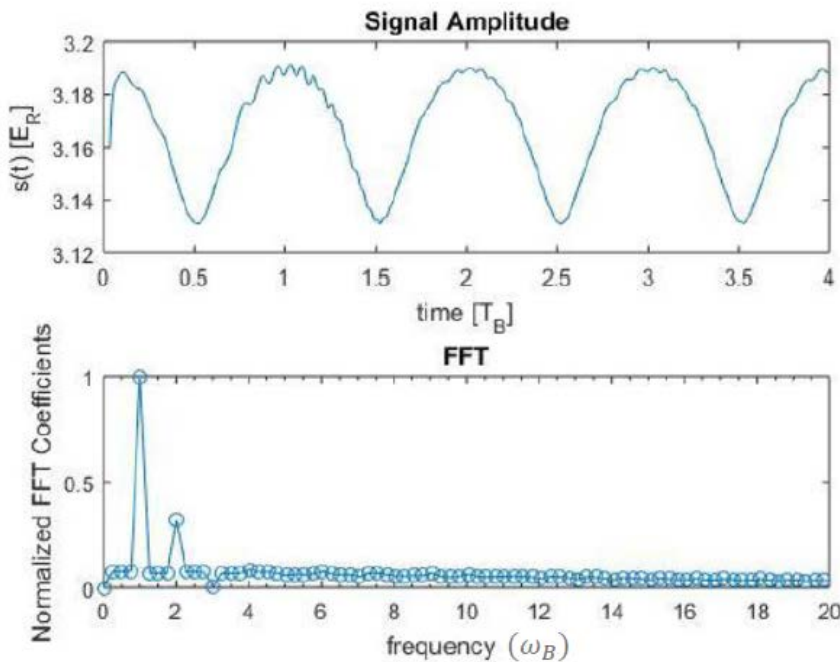
g = coupling (essentially vacuum Rabi frequency)

δ = detuning from cavity resonance

η = pumping rate

κ = cavity loss

Numerical simulations



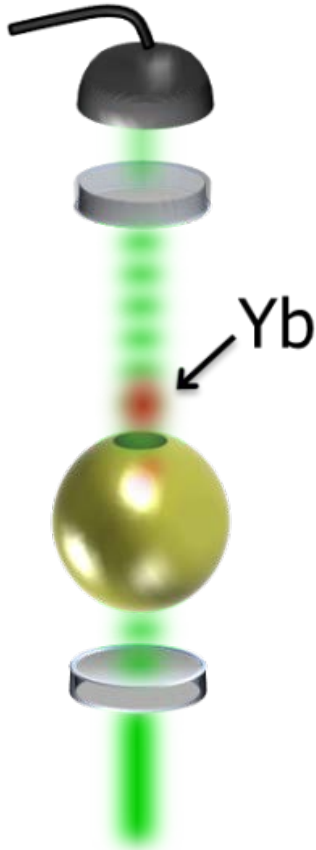
- 10^6 Yb atoms
- Bloch frequency **7.4 kHz**
- Cavity – **5 cm** long, **99.9%** reflectivity, **1 MHz** linewidth
- Lattice depth **$3 E_R$**
- Collective cooperativity **$>10,000$**

$$\rightarrow 10^{-7} \text{ g} / \sqrt{\text{Hz}}$$

Dark energy

Projected $10^{-9}g$ sensitivity in one day of integration

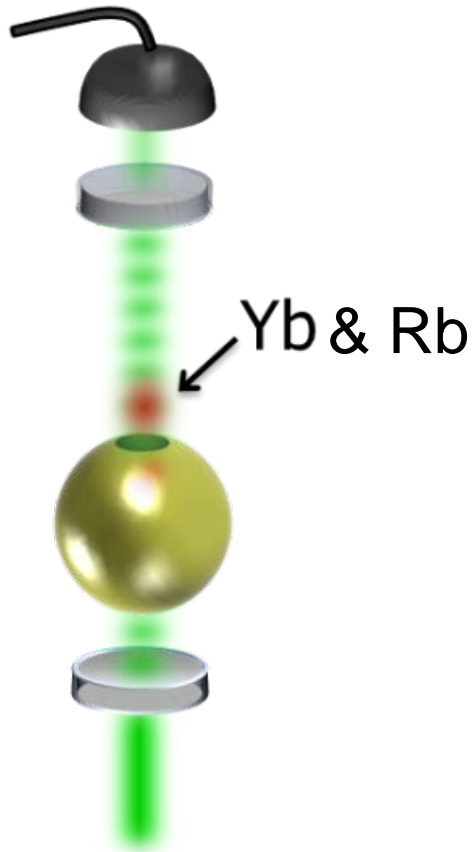
⇒ Rule out chameleons and constrain other scalar theories



Model	Description
Chameleon	Mass couples to matter density
Symmetron	Coupling depends on matter density
f(R) gravity	Equivalent to chameleon theory
Preferred scale	Maps to chameleon theory

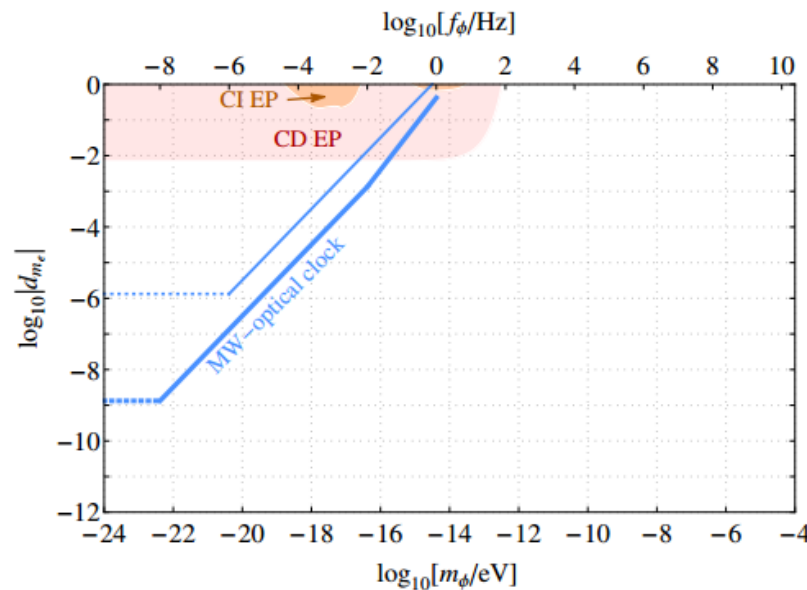
- Reduced vibration sensitivity / easier isolation
- Long coherence time

Dark matter

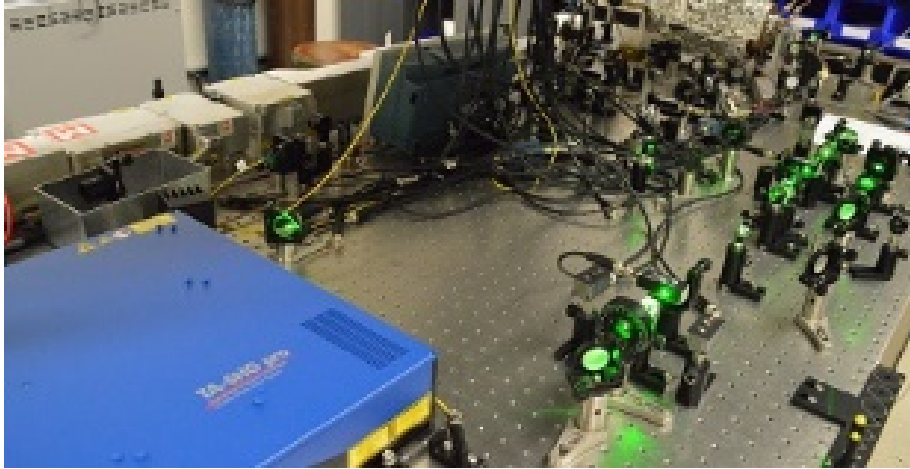


Time varying dilatons oscillate at Compton frequency.

10 kHz detection bandwidth for an EP test could improve constraints



Current status



Two stage cooling at 399 nm and 556 nm

Permanent magnet 2D MOT will be loaded from dispensers into 3D MOT

Zerodur cavity testing on benchtop

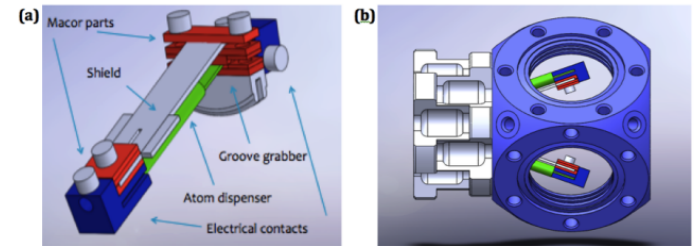
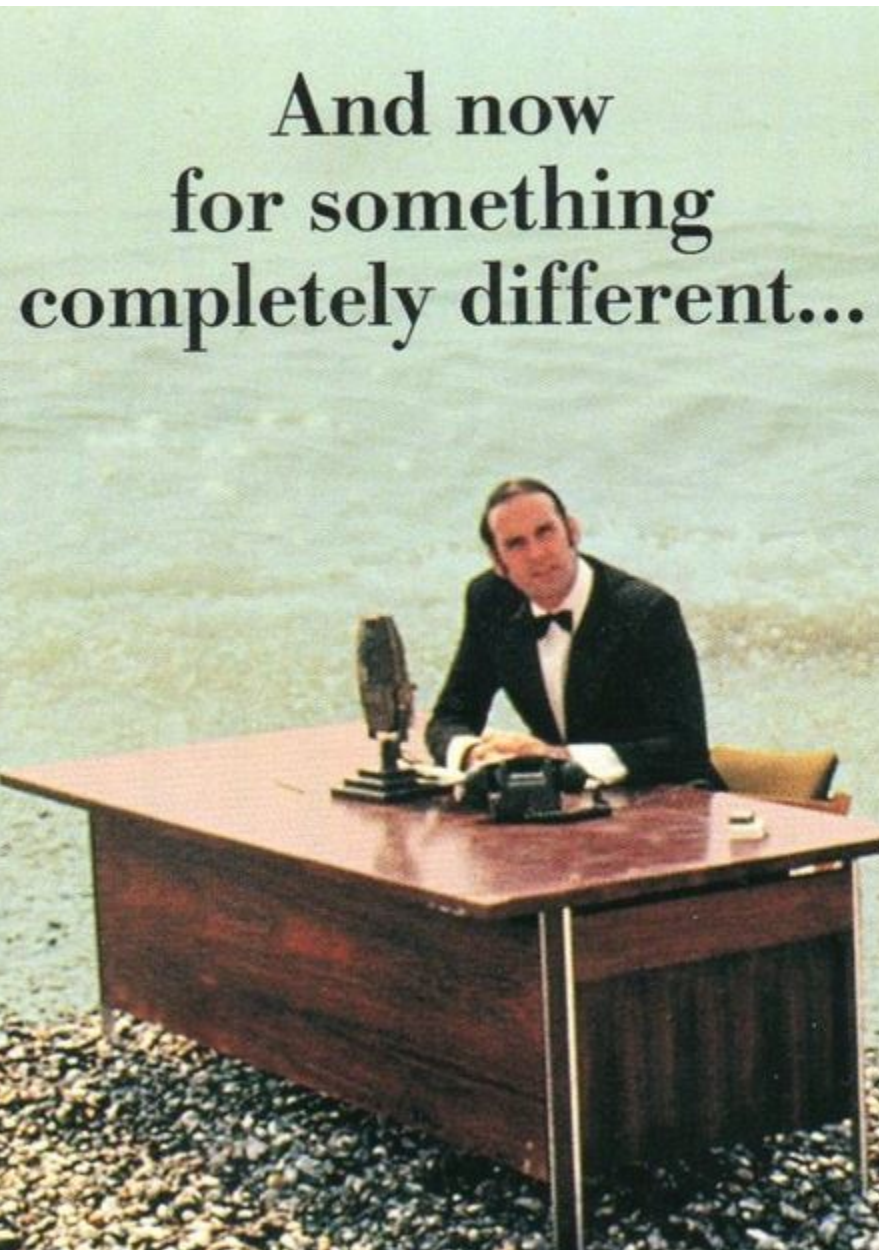


FIG. 3: (a) Atom dispenser assembly. For clarity, the atom dispenser is shaded green, the electrical contacts are blue, and the Macor insulating parts are red. (b) Atom dispensers mounted inside the 2D MOT chamber.



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HUNTER - sterile neutrino search





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HUNTER - Heavy Unseen Neutrinos by Total Energy-momentum Reconstruction



- Identifying the galactic dark matter is one of the key problems of modern physics.
- Many large scale experiments are searching for heavy “supersymmetric WIMPs” and will soon either find them or run into backgrounds.
- HUNTER will seek an entirely different dark matter candidate, the “sterile neutrino”, which would also fill a gaping hole in the Standard Model.

Jeff Martoff (Temple)
Eric Hudson (UCLA)
Paul Hamilton (UCLA)

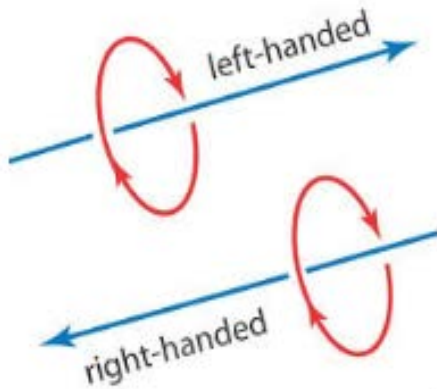
Peter F. Smith (UCLA)
Andrew Renshaw (Houston)
Hanguo Wang (UCLA)

Funding from W.M. Keck Foundation













Missing right-handed neutrino states

Quarks and Leptons have an elementary unit of spin

L H or R H relative to motion



Left handed	right handed
✓	✓
✓	✓
✓	✓
✓	MISSING

$\approx 2.3 \text{ MeV}/c^2$ $2/3$ $1/2$  up	$\approx 1.275 \text{ GeV}/c^2$ $2/3$ $1/2$  charm	$\approx 173.07 \text{ GeV}/c^2$ $2/3$ $1/2$  top
$\approx 4.8 \text{ MeV}/c^2$ $-1/3$ $1/2$  down	$\approx 95 \text{ MeV}/c^2$ $-1/3$ $1/2$  strange	$\approx 4.18 \text{ GeV}/c^2$ $-1/3$ $1/2$  bottom
$0.511 \text{ MeV}/c^2$ -1 $1/2$  electron	$105.7 \text{ MeV}/c^2$ -1 $1/2$  muon	$1.777 \text{ GeV}/c^2$ -1 $1/2$  tau
$< 2.2 \text{ eV}/c^2$ 0 $1/2$  electron neutrino	$< 0.17 \text{ MeV}/c^2$ 0 $1/2$  muon neutrino	$< 15.5 \text{ MeV}/c^2$ 0 $1/2$  tau neutrino

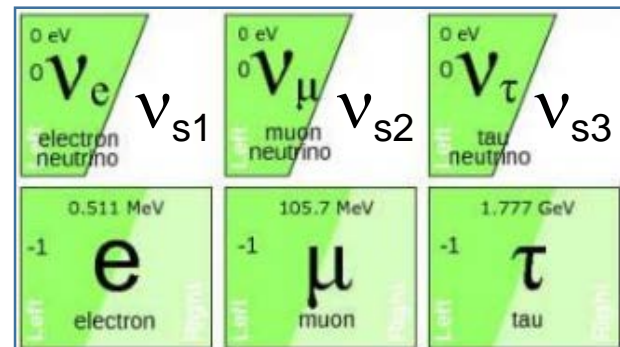
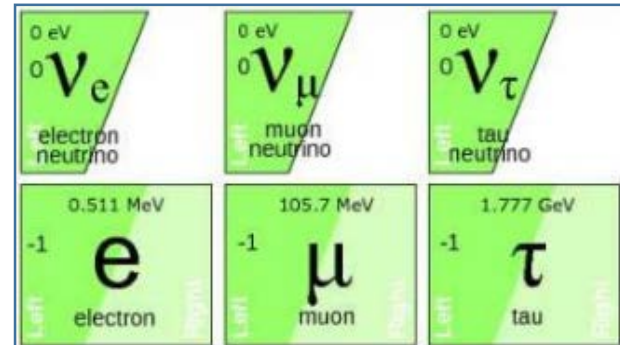
Unseen in existing neutrino interactions

Addition of sterile neutrino states

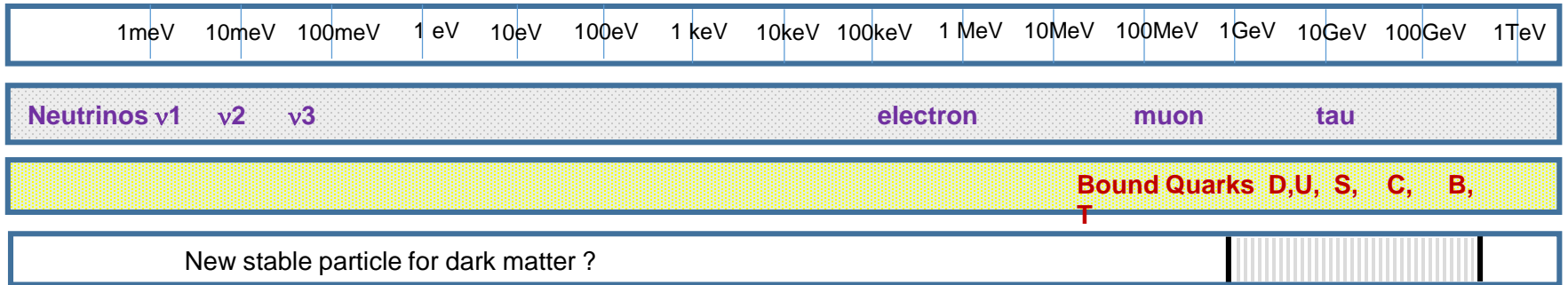
- Pre-1980: neutrinos thought zero mass with only LH states (& anti-neutrinos RH)
- Post-1980: Neutrinos confirmed to have small but non-zero masses
- **Non-zero mass indicates both LH and RH states should exist, but not yet seen**



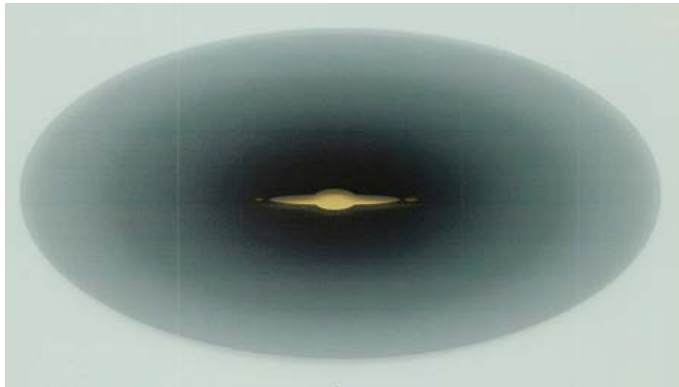
- Thus the missing neutrinos must have either high mass or Interaction strength much weaker than known neutrinos
- Hence named '**sterile**' (= '**quasi sterile**')
- **No clear prediction of mass**



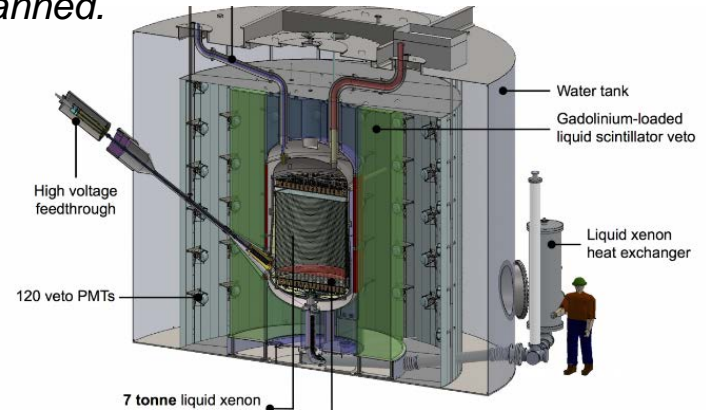
Mass range of the known 'elementary' particles



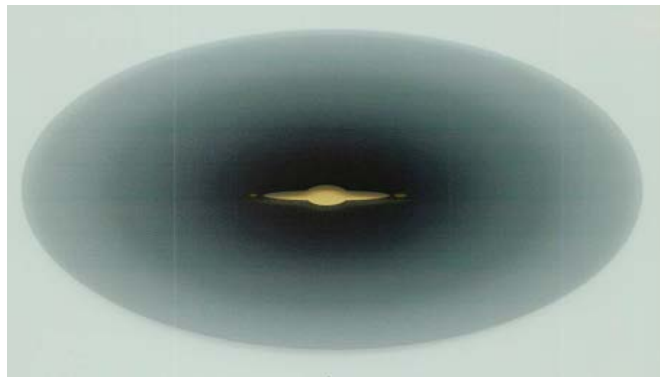
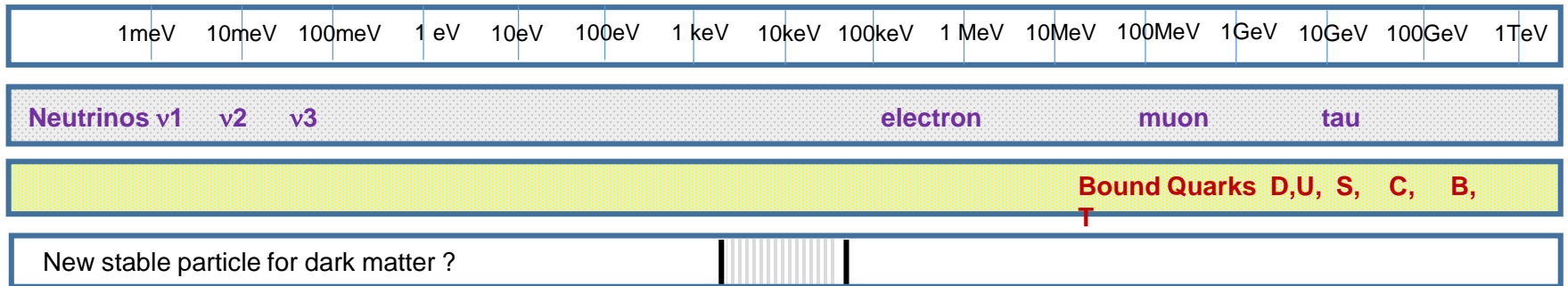
Galactic dark matter Problem



No heavy particle dark matter signals seen in underground ton-scale detectors or at Large Hadron Collider. Multi-ton detectors planned.



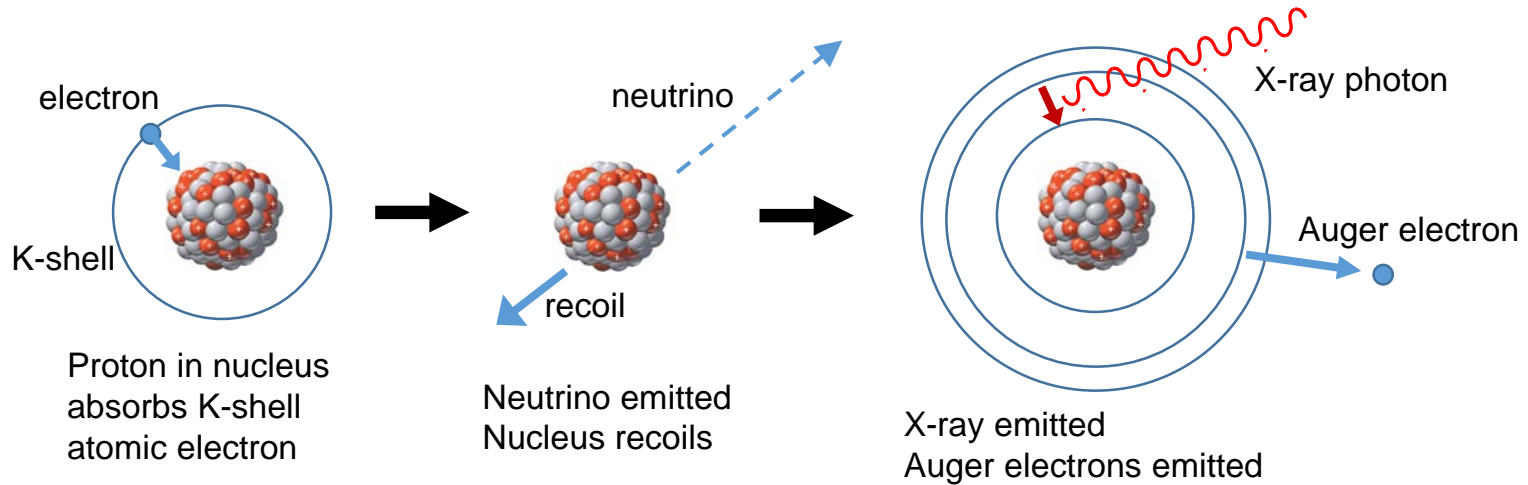
Galactic Dark Matter - keV-mass sterile neutrinos can provide the answer



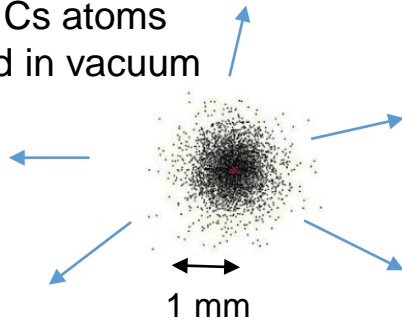
Theoretical calculations show

- ~ 10 keV sterile neutrino mass range fits known galactic dark matter density.
- Weak relative coupling $< 10^{-4}$ gives stability for universe lifetime.
- Direct detection not feasible in the foreseeable future

Proposed K-capture experiment: measuring the mass of an unseen neutrino



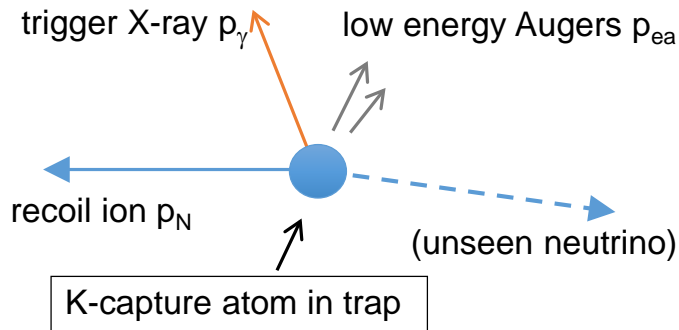
- Cloud of ^{131}Cs atoms suspended in vacuum



- measure momentum of ion, X-ray, e, to calculate neutrino momentum and mass
- **This will find rare keV-mass sterile neutrinos up to Q value of decay (350 keV for ^{131}Cs)**
- **Fraction of signal events gives relative coupling**

Summary of HUNTER principle (Heavy Unseen Neutrinos by Total Energy-momentum Reconstruction)

Measurements required:

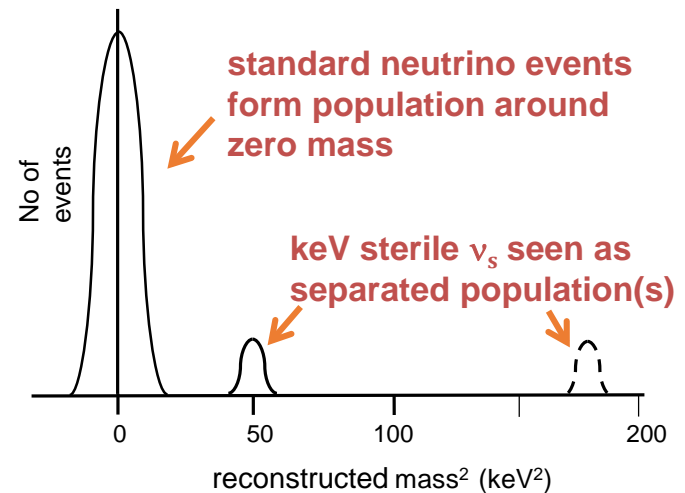


Mass reconstruction formula:

$$m_\nu^2 = [Q - E_a - E_\gamma - E_N]^2 - [\mathbf{p}_\gamma + \mathbf{p}_{ea} + \mathbf{p}_N]^2$$

missing energy missing momentum

Reconstructed mass spectra:



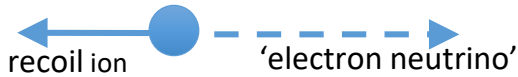
- This is the only known method of giving a separated population of sterile neutrino events
- Can find sterile neutrinos independently of whether they form all or part of the dark matter

see Peter F. Smith, arxiv:1607.06876 for details

How can the electron neutrino from K-capture become a sterile neutrino ?

neutrino flavors ($\nu_e, \nu_\mu, \nu_\tau, \nu_{s1}, \nu_{s2}, \nu_{s3}$) \leftrightarrow mixtures of definite mass ($\nu_1, \nu_2, \nu_3, \nu_4, \nu_5, \nu_6$)

$$\begin{array}{c}
 \nu_e \\
 \nu_\mu \\
 \nu_\tau \\
 \nu_{s1} \\
 \nu_{s2} \\
 \nu_{s3}
 \end{array}
 =
 \begin{array}{c}
 c_{11} \ c_{12} \ c_{13} \ c_{14} \ c_{15} \ c_{16} \\
 c_{21} \ c_{22} \ c_{23} \ . \ . \ . \\
 c_{31} \ c_{43} \ c_{33} \ . \ . \ . \\
 . \ . \ . \ . \ . \ . \\
 . \ . \ . \ . \ . \ . \\
 . \ . \ . \ . \ . \ .
 \end{array}
 \begin{array}{c}
 \nu_1 \\
 \nu_2 \\
 \nu_3 \\
 \nu_4 \\
 \nu_5 \\
 \nu_6
 \end{array}
 \text{ mass eigenstates}$$



$$\nu_e = c_{11}\nu_1 + c_{12}\nu_2 + c_{13}\nu_3 + c_{14}\nu_4 + c_{15}\nu_5 + c_{16}\nu_6$$

standard neutrino mass states

admixture of 'sterile' neutrinos

Lower probability but will be seen as rare events

131-Cs Decay

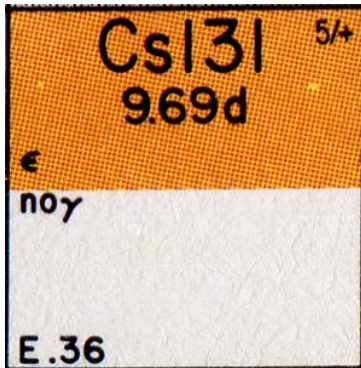
$t_{1/2}=9.7$ da, EC 100%,
 $Q_{EC}=355$ keV

131-Cs \rightarrow 131-Xe (stable)

+ ν

+ x-rays (4-35 keV)

+ Auger e-'s (3-150 eV)



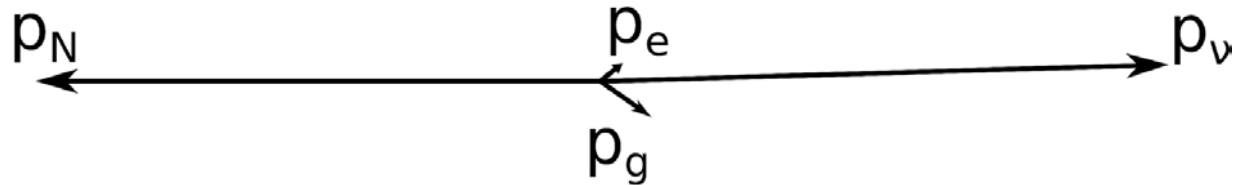
No penetrating radiation, no radioactive daughter.

Commercially available

(IsoRay brachytherapy seeds)

\$10K/order + \$1K/Ci

Basic Kinematics



$$p_N, p_v \sim .35 \text{ MeV}/c \quad v_N \sim 800 \text{ m/s}$$

$$p_e \sim .01 \text{ MeV}/c, p_g \sim .03 \text{ MeV}/c$$

$$m_v^2 = [Q - E_e - E_g - E_N]^2 - [\mathbf{p}_g + \mathbf{p}_e + \mathbf{p}_N]^2$$

Ignore e and x-ray, calculate effect of m_v on p_N :

$$p^2/2m_N + (p^2 + m_v^2)^{1/2} - m_v = Q$$

Accurate first-order solution: $p = Q(1 - m_v^2/2Q^2)$

For $m_v = 10 \text{ keV}$, effect is .04% of p !

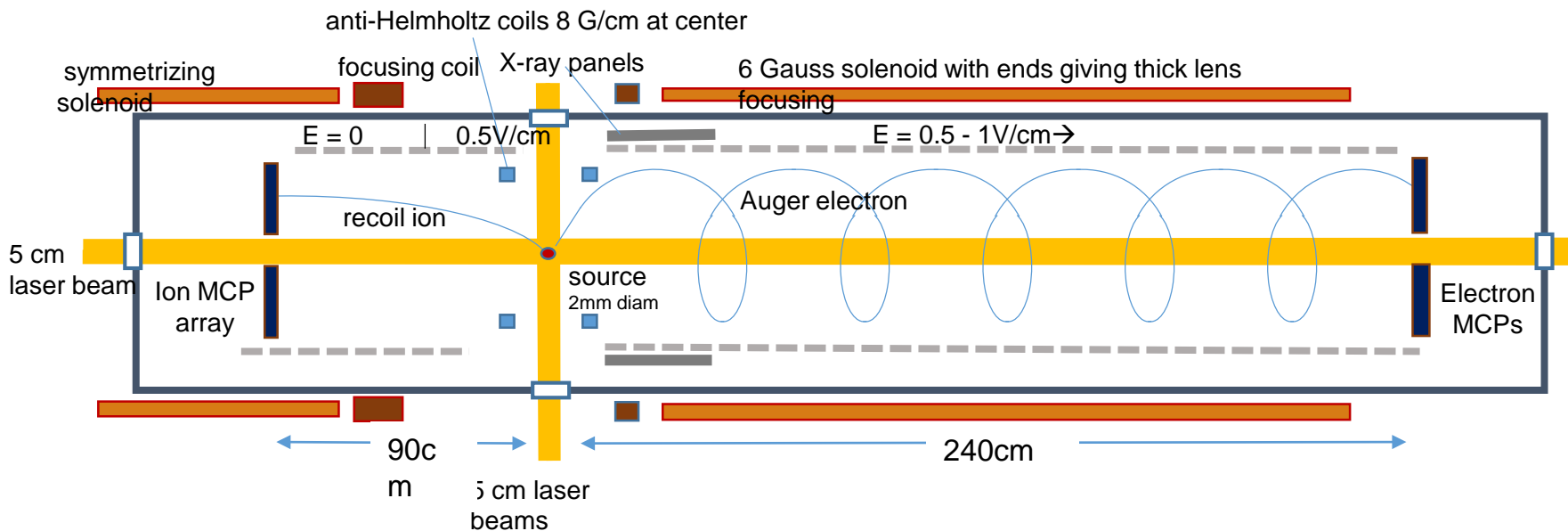
This sets the scale of measurement accuracy needed.

Note: this δp is equal to the thermal momentum p_{th} at $150 \mu\text{K}$.

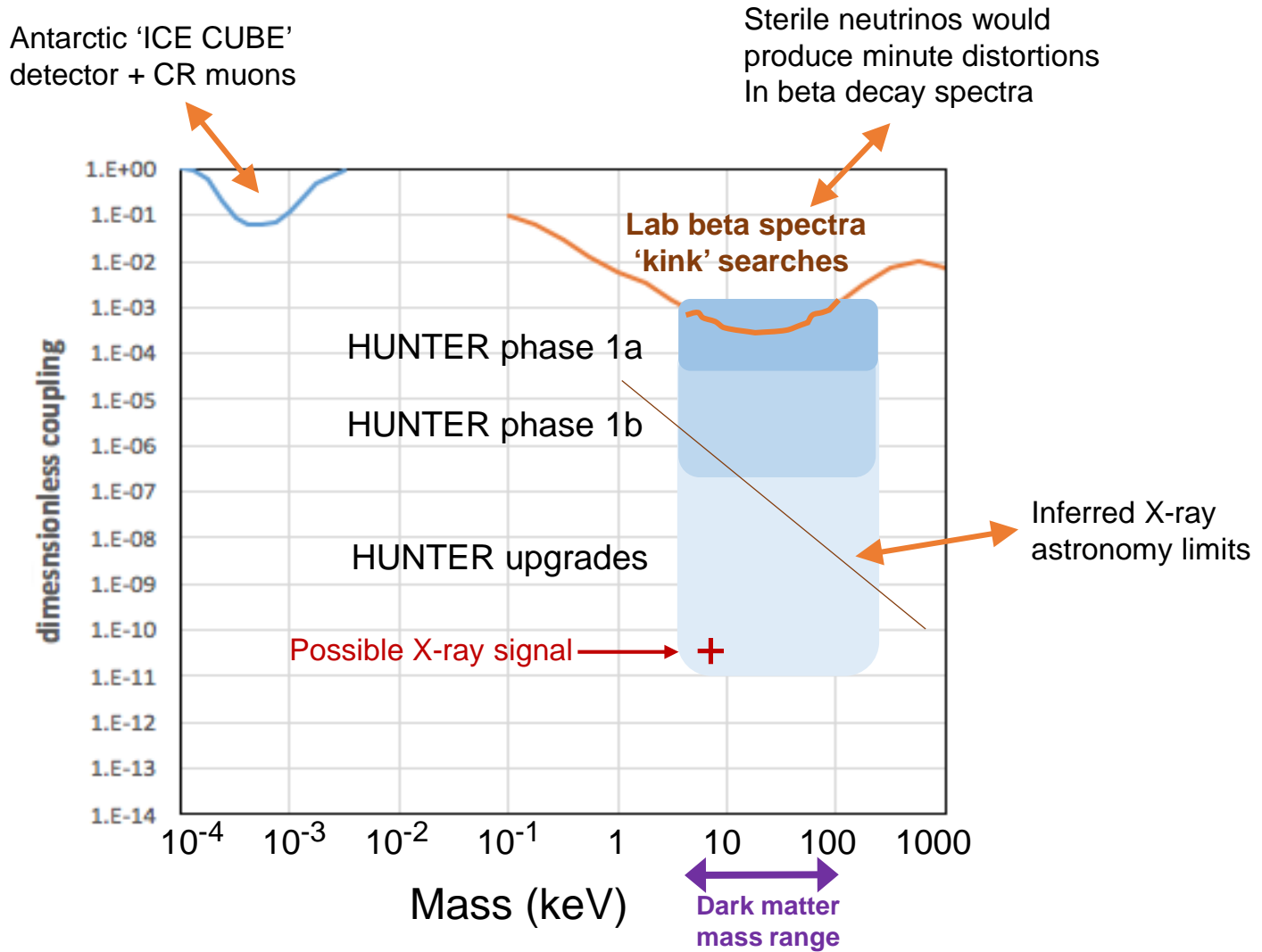
Trapped atoms must be colder than this by factor 3 or more.

Practical configuration

- 4m long vacuum vessel
- Source trapped in intersecting laser beams
- X-ray photon detected by scintillator + multi-anode SiPM array
- Recoil ion and Auger electrons directed by electric fields to MCP arrays
- Longer electron path length confined by magnetic spiraling



Existing limits and future coverage of HUNTER experiment

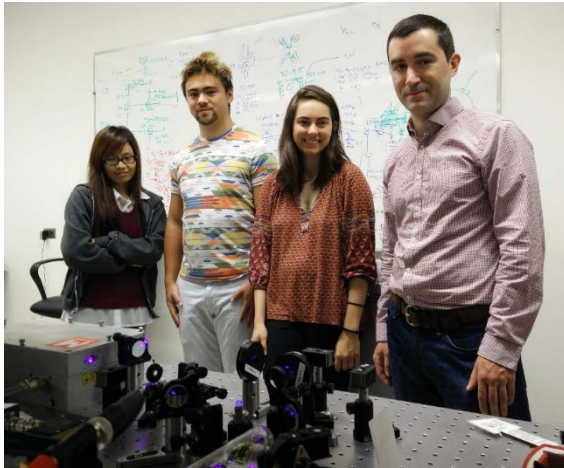


- **Matter wave interferometry**
- **Test of dark energy at Berkeley**
 - 100x improvement on first limits of screened scalar fields
- **Experiments in development at UCLA**
 - Cavity detection of Bloch oscillations for continuous force sensing
 - HUNTER – sterile neutrino search
 - Precision ion interferometry



UCLA

Thanks



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